



# Nuclear liquid-gas phase transition with machine learning

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6.28.2022 Journal club



PHYSICAL REVIEW RESEARCH **2**, 043202 (2020)

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## Nuclear liquid-gas phase transition with machine learning

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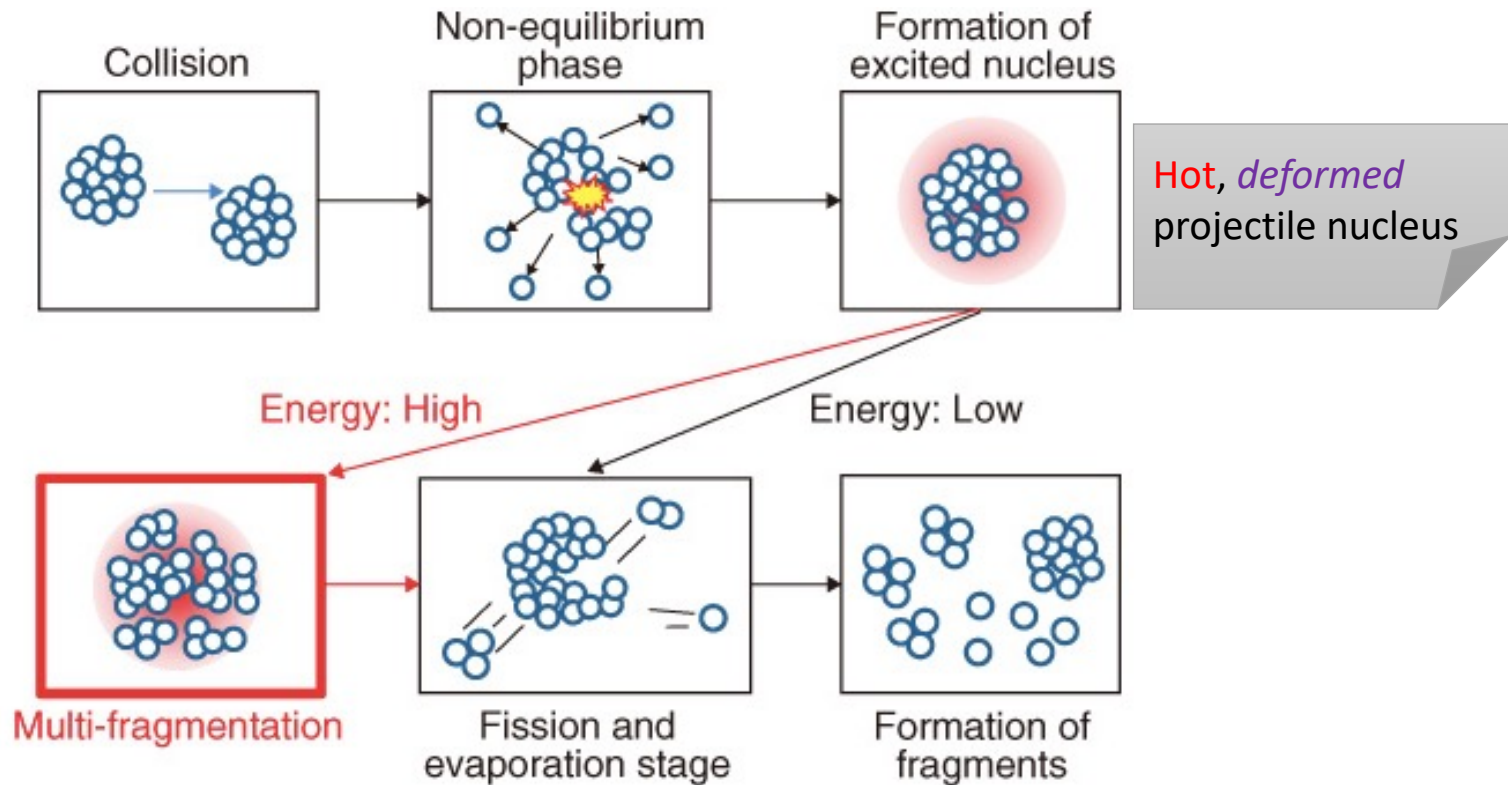
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1. Preliminaries
2. Experimental data processing
3. Machine learning results
4. Highlights of this work

# Preliminaries

## Nuclear multifragmentation $^{40}\text{Ar}$ on $^{27}\text{Al}$ & $^{48}\text{Ti}$ at 47 MeV/nucleon

@TAMU K500 superconducting cyclotron  
with detector NIMROD

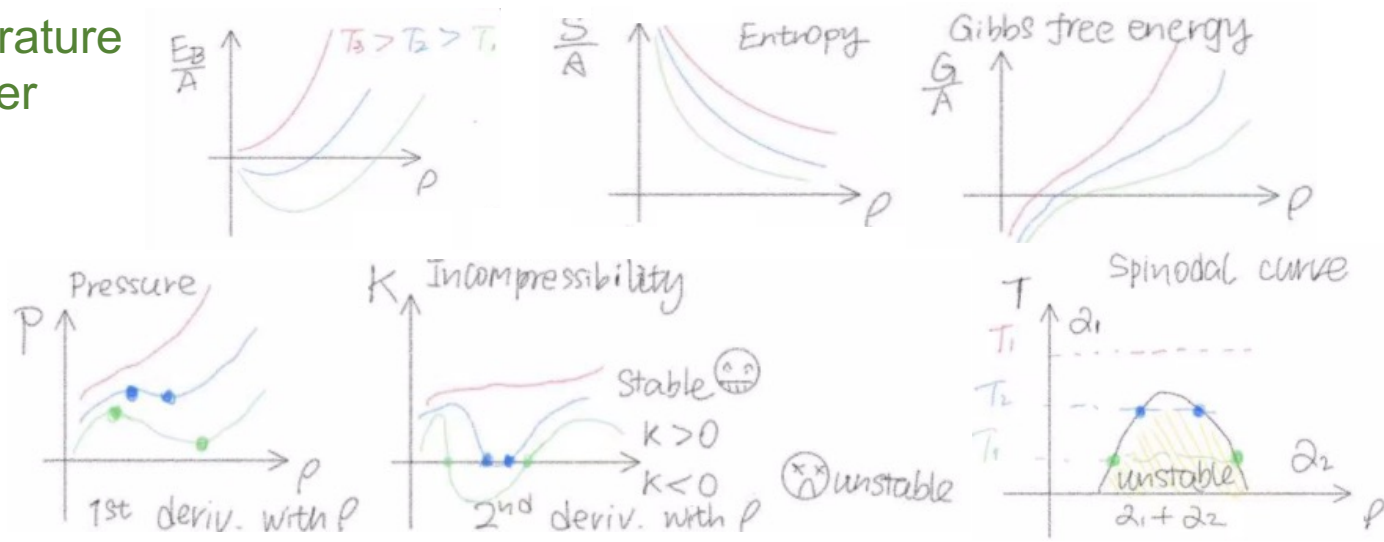


See more on [https://rdreview.jaea.go.jp/review\\_en/2013/e2013\\_8\\_2.html](https://rdreview.jaea.go.jp/review_en/2013/e2013_8_2.html)

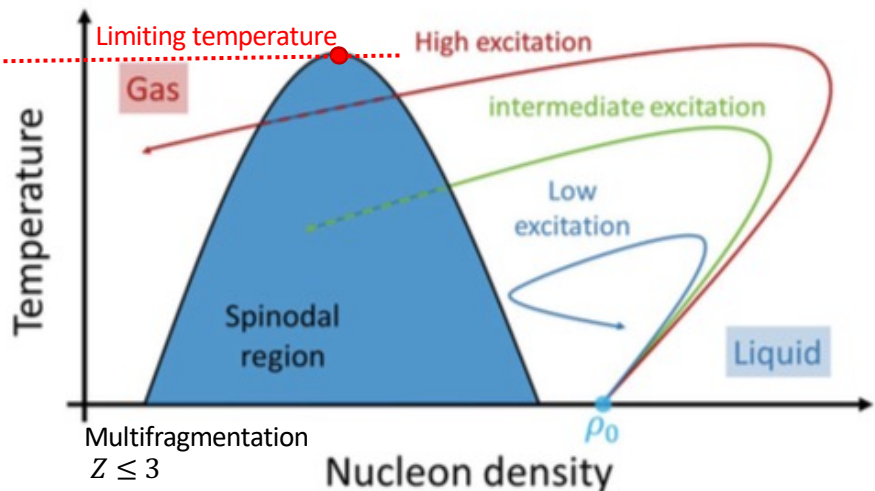
# Spinodal decomposition



## Finite-temperature nuclear matter



## Heavy-ion reaction



**Liquid like**  
low excitation (0.9–2.8 MeV)

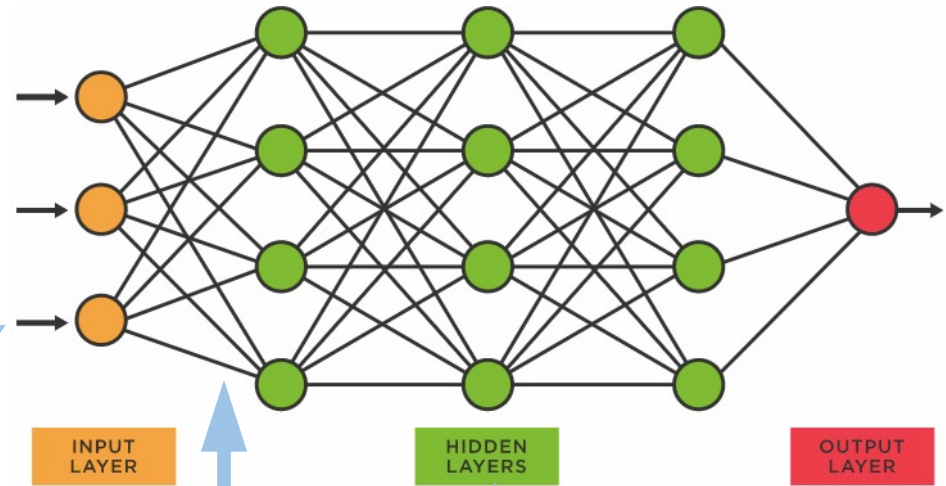
**Coexistence**  
intermediate excitation (5.3–5.4 MeV)

**Gas like**  
high excitation (8.1–13.0 MeV)

# Machine learning: neural network (NN)



A video nicely introduces NN: <https://www.youtube.com/watch?v=aircAravnKk>



Raw data  
(preprocessing)

training set

$$\mathbf{X} = \{\mathbf{x}_1 \dots \mathbf{x}_N\}$$

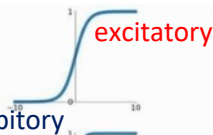
testing set

Activation functions

$$f(\mathbf{w}_i \mathbf{x}_i + b_i)$$

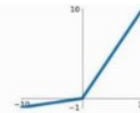
**Sigmoid**

$$\sigma(x) = \frac{1}{1+e^{-x}}$$



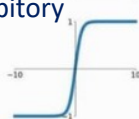
**Leaky ReLU**

$$\max(0.1x, x)$$



**tanh**

$$\tanh(x)$$



**Maxout**

$$\max(w_1^T x + b_1, w_2^T x + b_2)$$

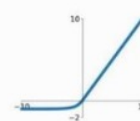
**ReLU**

$$\max(0, x)$$



**ELU**

$$\begin{cases} x & x \geq 0 \\ \alpha(e^x - 1) & x < 0 \end{cases}$$



Hidden layers

$$y_i = g(\mathbf{x}_i) \\ \approx f_\ell(\dots f_1(\mathbf{W}_1 \mathbf{x}_i + \mathbf{b}_1))$$

Training (supervised)

finding  $\{\mathbf{W}_\ell, \mathbf{b}_\ell\}$  to minimize object function

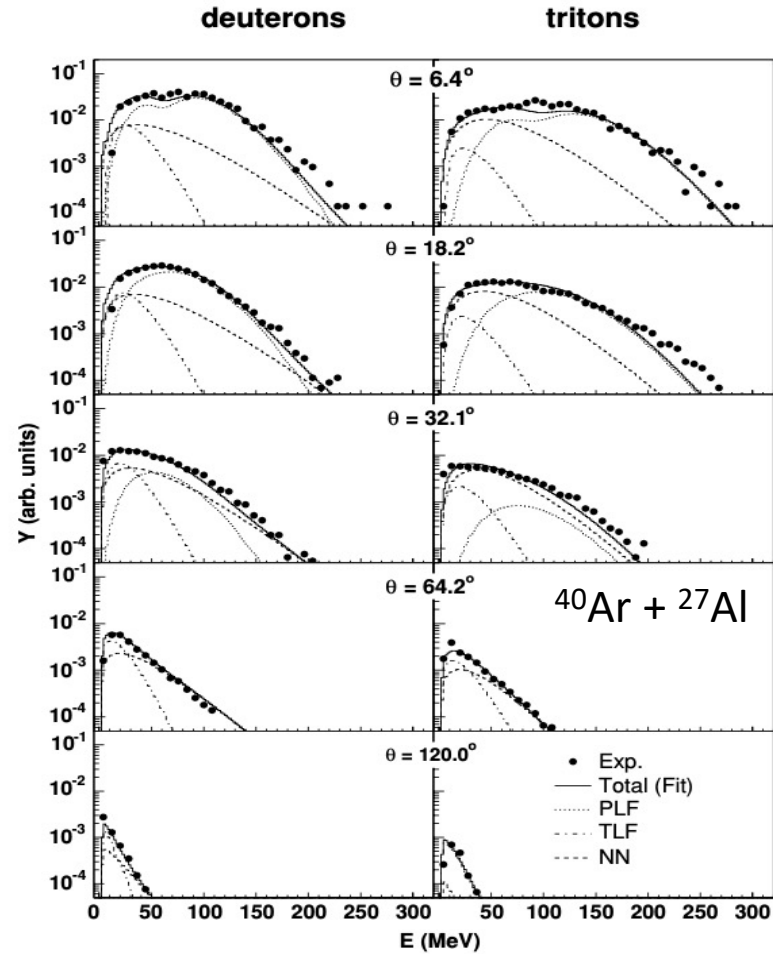
# Experimental data pre-processing

Y. G. Ma, et al., PhysRevC71, 054606 (2005)

It is well known that the laboratory frame kinetic energy spectra of most light ejectiles can be reproduced with the assumption of emission from three different sources: (i) a quasiprojectile (QP) source, (ii) an intermediate velocity or nucleon-nucleon (NN) source, and (iii) a quasitarget (QT) source. To better understand origins of the emitted particles the ideal situation would be to have the ability to attribute each particle to its source on an event-by-event basis. However the spectral distributions from the different sources overlap significantly, making such an attribution not possible. Previous techniques to reconstruct QP have included identifying high-velocity components of the QP [54,55] or treating only particles emitted in the forward hemisphere in the projectile frame and then assuming identical properties for particles emitted in the backward hemisphere to recreate the QP source [33,34]. Such a technique is limited in its application and not suited to situations in which fluctuations are to be investigated.

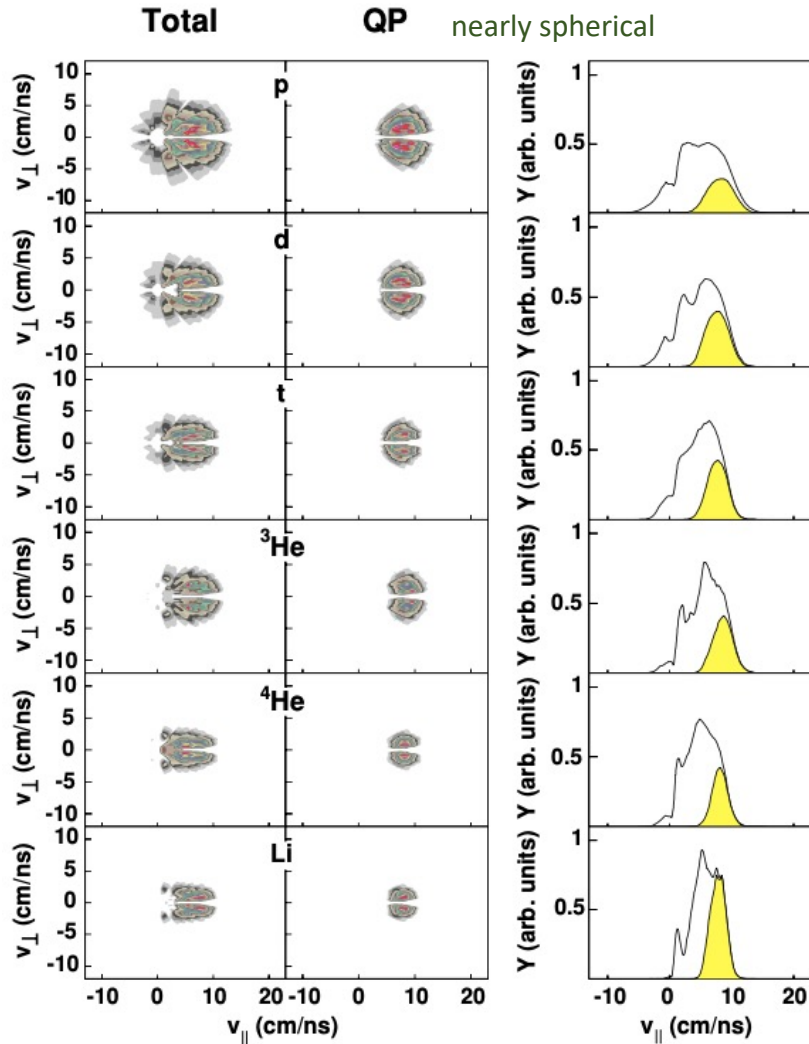
For this work we have developed a new method for the assignment of each light-charged particle (LCP) to an emission source. This is done with a combination of three source fits and Monte Carlo sampling techniques. We first obtain the laboratory energy spectra for different LCP at different laboratory angles and reproduce them using the three source fits. In the laboratory frame, the energy spectra of LCP can be modeled as the overlap of emission from three independent moving equilibrated sources (i.e., the QP, NN, and QT sources). For

## Reconstructed differential cross section through **three-source fitting**

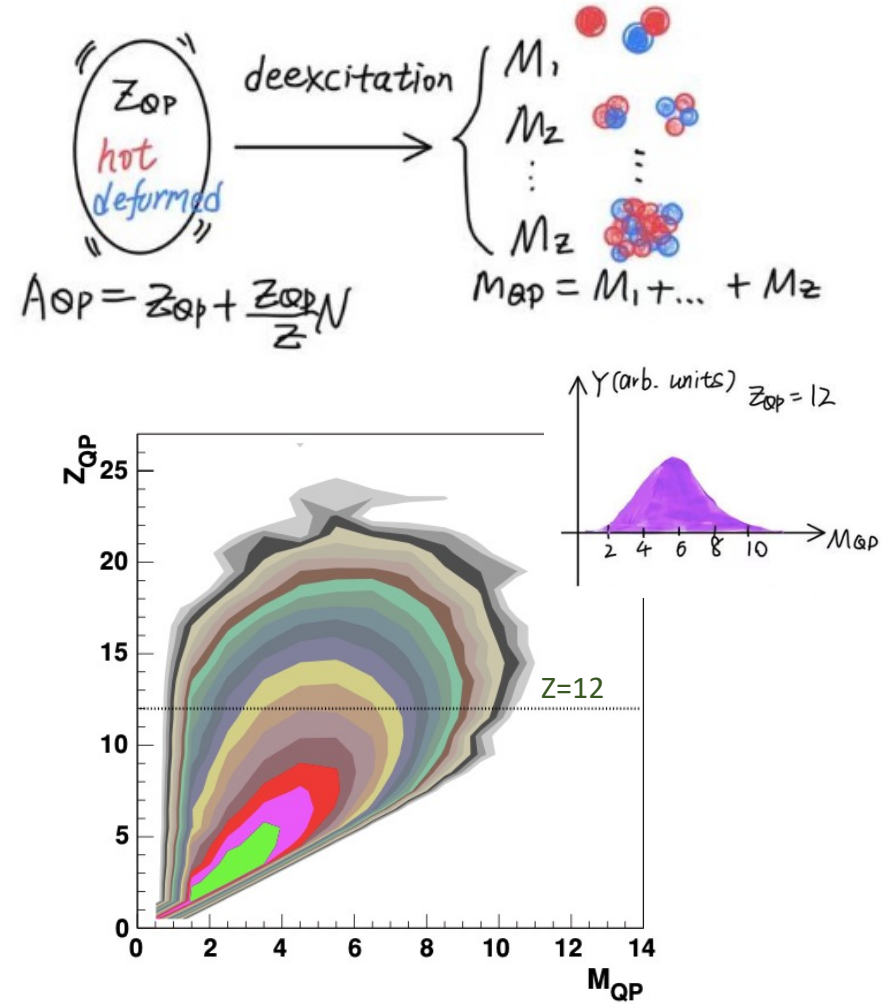




## Velocity distributions



## QP deexcitation



Event-by-event reconstructed QP

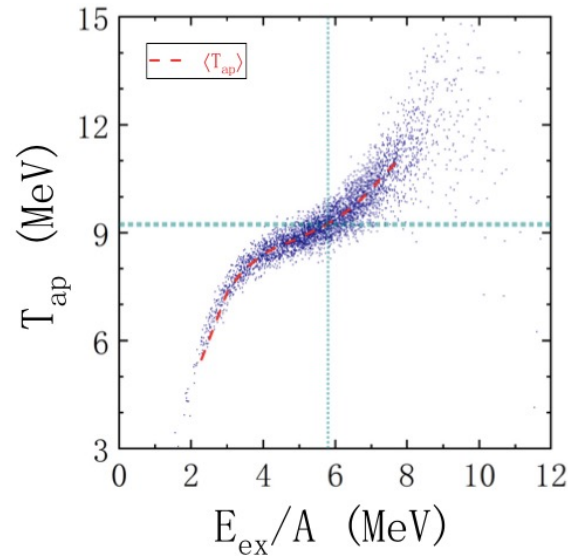
$$\mathbf{x}_i = \{M_1, \dots, M_{12}, E_{\text{ex}}/A, \langle T_{\text{ap}} \rangle\}$$

The caloric curve

$$E_{\text{ex}} = \sum_{i=1}^{M_{\text{QP}}} E_i^{\text{kin}} + \frac{3}{2} M_n T - Q$$

$$\langle T_{\text{ap}} \rangle = \sqrt{\frac{\langle Q_{xy}^2 \rangle - \langle Q_{xy} \rangle^2}{4m^2}}$$

$$Q_{xy} = p_x^2 - p_y^2$$



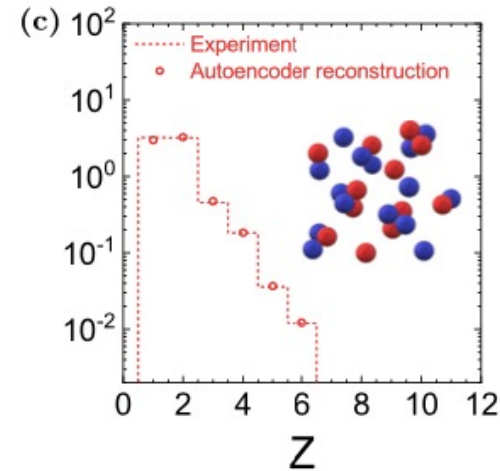
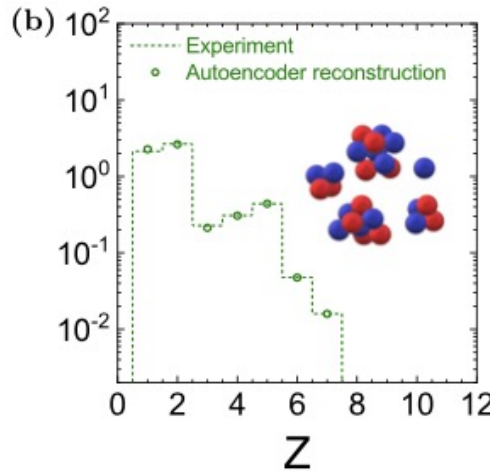
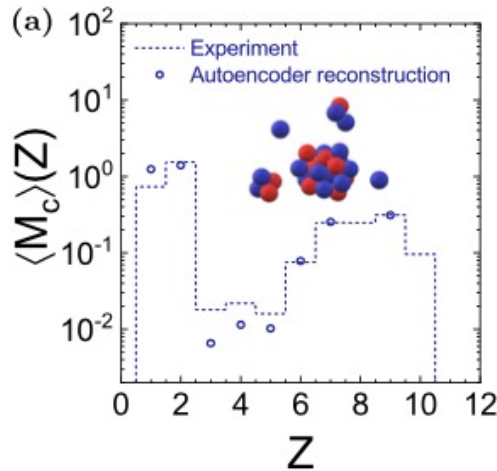
$$\tilde{c} \equiv \frac{d(E_{\text{ex}}/A)}{dT_{\text{ap}}}$$



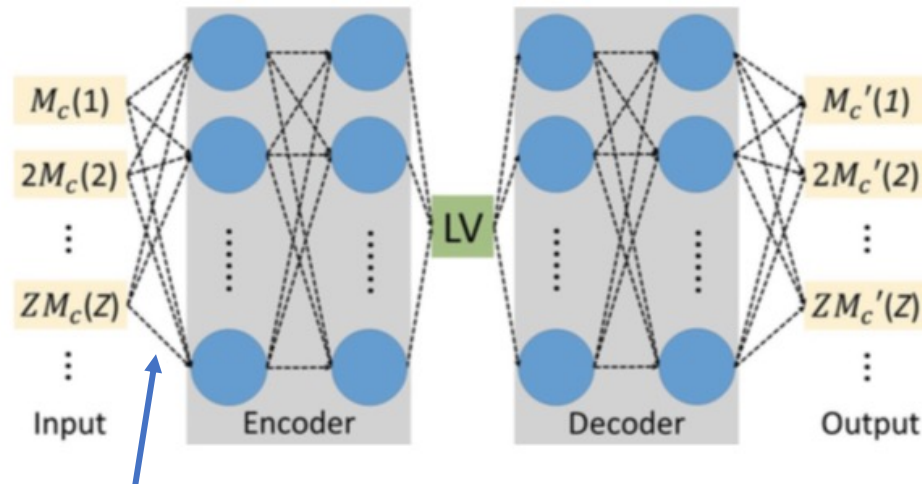
$$E_{\text{ex}}/A = 5.6 \pm 0.5 \text{ MeV}$$

$$T_{\text{ap}} = 8.3 \pm 0.5 \text{ MeV}$$

The averaged charge multiplicity distribution  $\langle M_c \rangle(Z)$  of the QP fragments



## Machine learning results



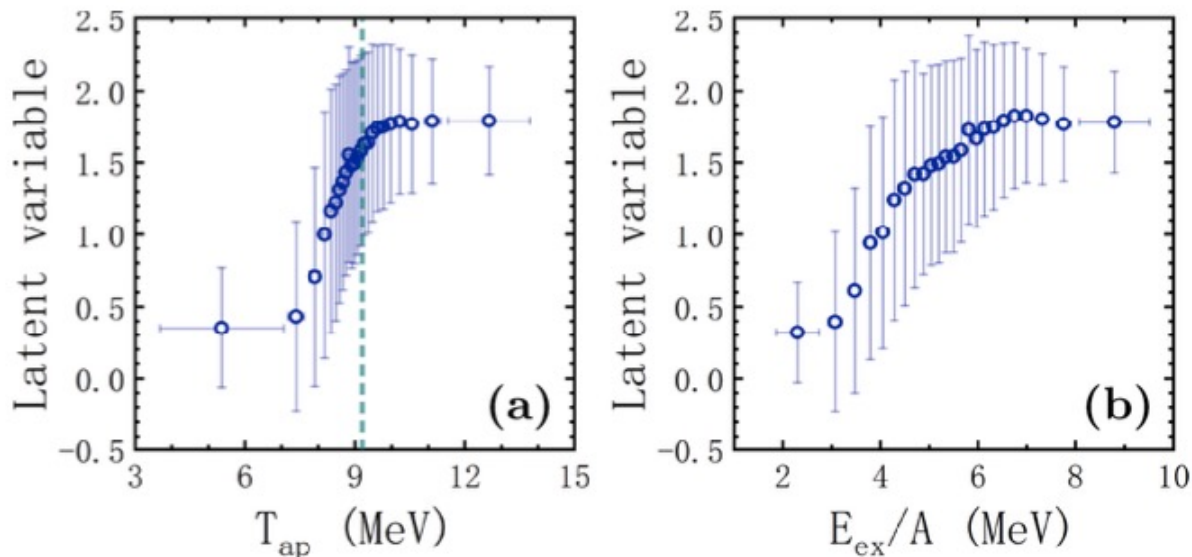
Activation function  $ReLU(\mathbf{W}\mathbf{x} + \mathbf{b})$

Layer	Neuron number	Activation
Input	12	$\tanh$
Encoder first layer	8-64	$\tanh$
Encoder second layer	4-32	$\tanh$ or $ReLU$
Latent variable	1	$\tanh$
Decoder first layer	4-32	$\tanh$
Decoder second layer	8-64	$ReLU$
Output	12	-

Loss function with regulator to be minimized

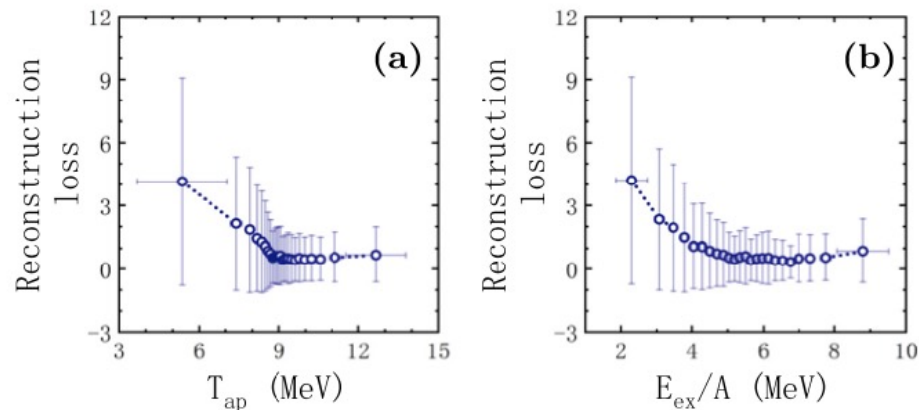
$$\tilde{C}(\tilde{\mathbf{y}}, \mathbf{y}) = (\tilde{\mathbf{y}} - \mathbf{y})^2 + l_2(\|\mathbf{W}\|^2/2 + \|\mathbf{b}\|^2/2).$$

The mean and standard deviation of the latent variable in different  $T_{ap}$  and  $E_{ex}/A$  bins



The mean and standard deviation of the event-by-event reconstruction loss of the autoencoder network

$$\sum_{Z=1}^{12} |ZM_c(Z) - ZM'_c(Z)|$$



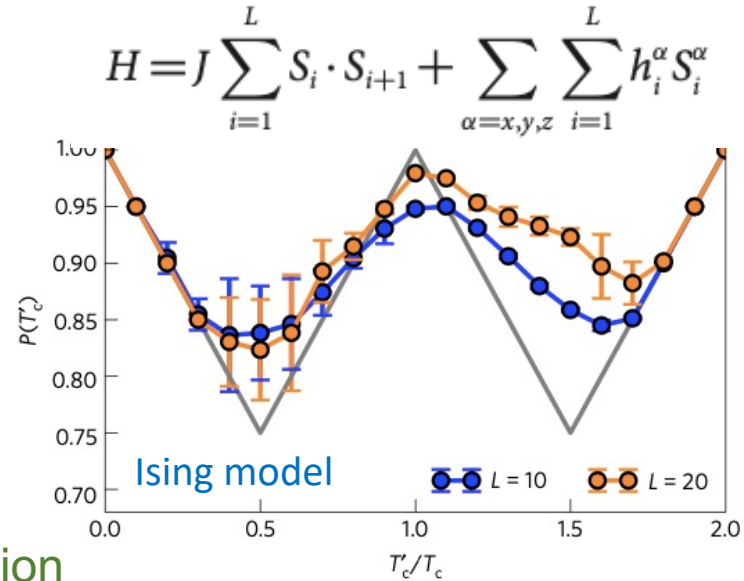
# 'Confuse' a Bayesian neural network



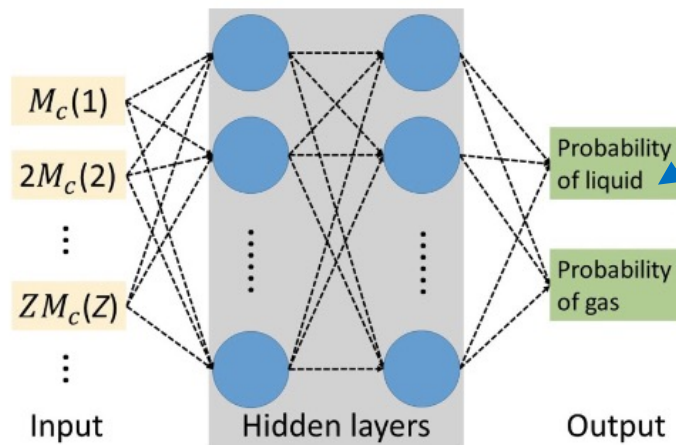
## Confusion scheme

[E P L van Nieuwenburg, et al., NatPhys13, 435-439 \(2017\)](#)

In this Letter we demonstrate that it is possible to find the correct labels, by *purposefully mislabelling the data* and evaluating the performance of the machine learner.



## Confusion scheme for $T'_{ap}$ in heavy-ion reaction

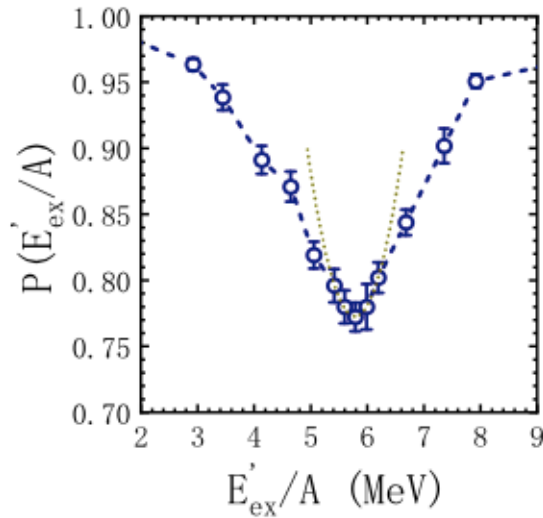
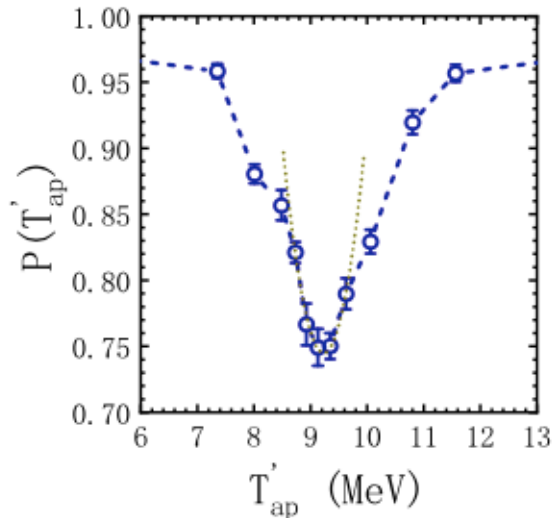


Give a wrong  $T'_{ap}$  for classification & train the network

# 'Confuse' a Bayesian neural network

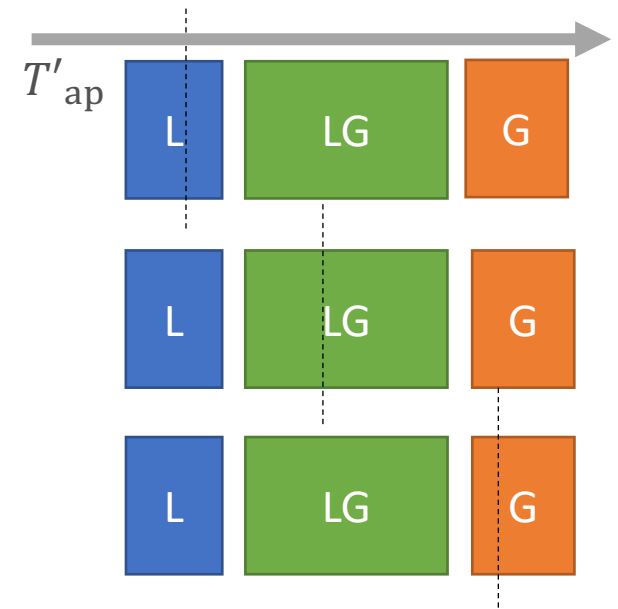


The performance in the test data set



$$E_{ex}/A = 5.79 \pm 0.02 \text{ MeV}$$

$$T_{ap} = 9.24 \pm 0.04 \text{ MeV}$$

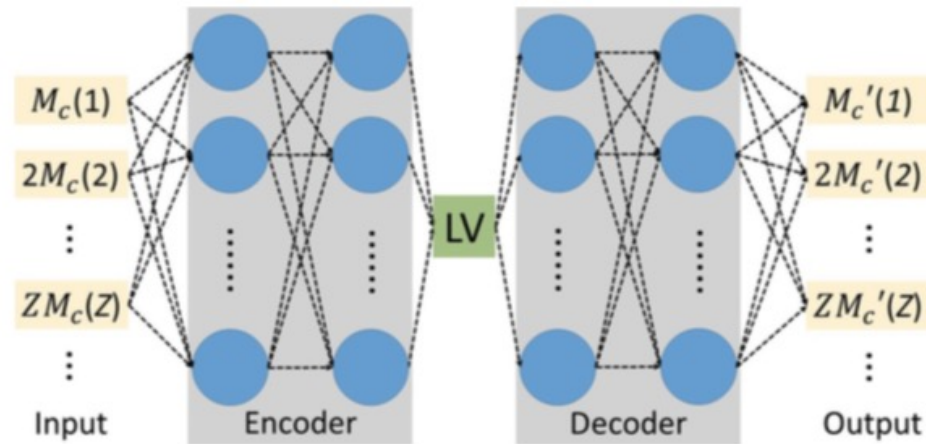


## Highlights of this work

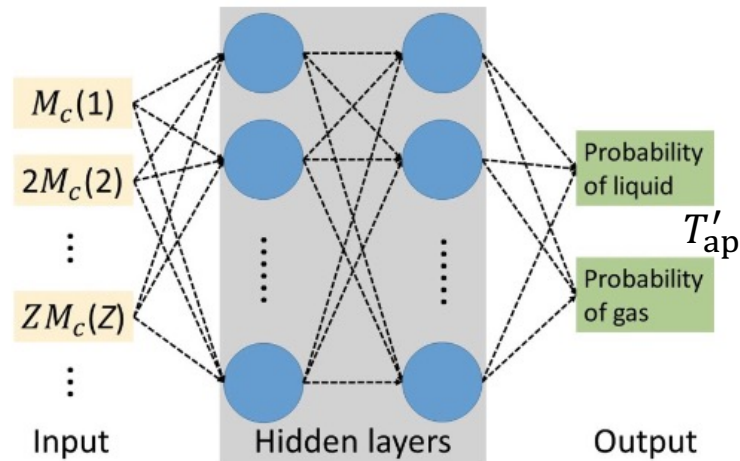


## Autoencoder & latent variable

LV as kind of order parameter



## Confusion scheme



**Thank you for your attention**