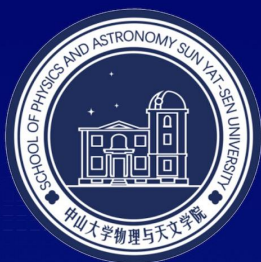


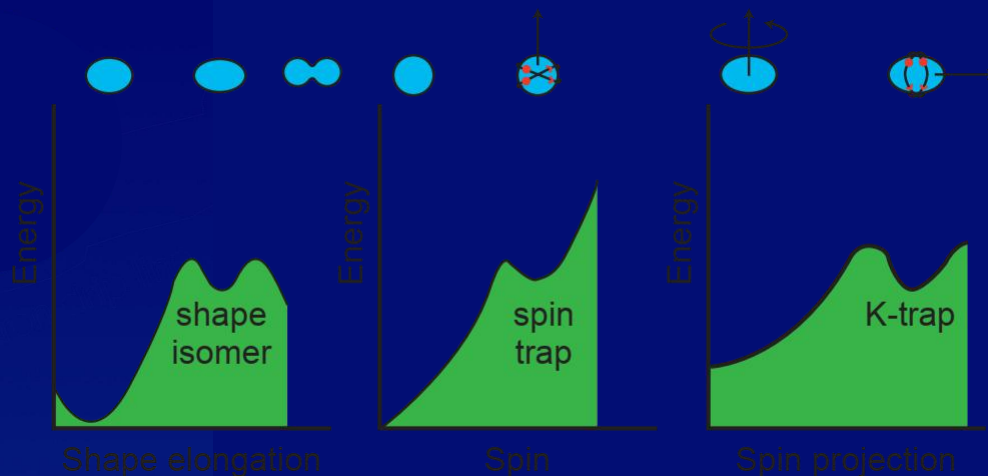
# Metastable States of $^{92,94}\text{Se}$ : Identification of an Oblate $K$ Isomer of $^{94}\text{Se}$ and the Ground State Shape Transition between $N = 58$ and $60$

C. Lizarazo et.al.,

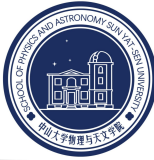


Changfeng Jiao (焦长峰)

School of Physics and Astronomy  
Sun Yat-sen University



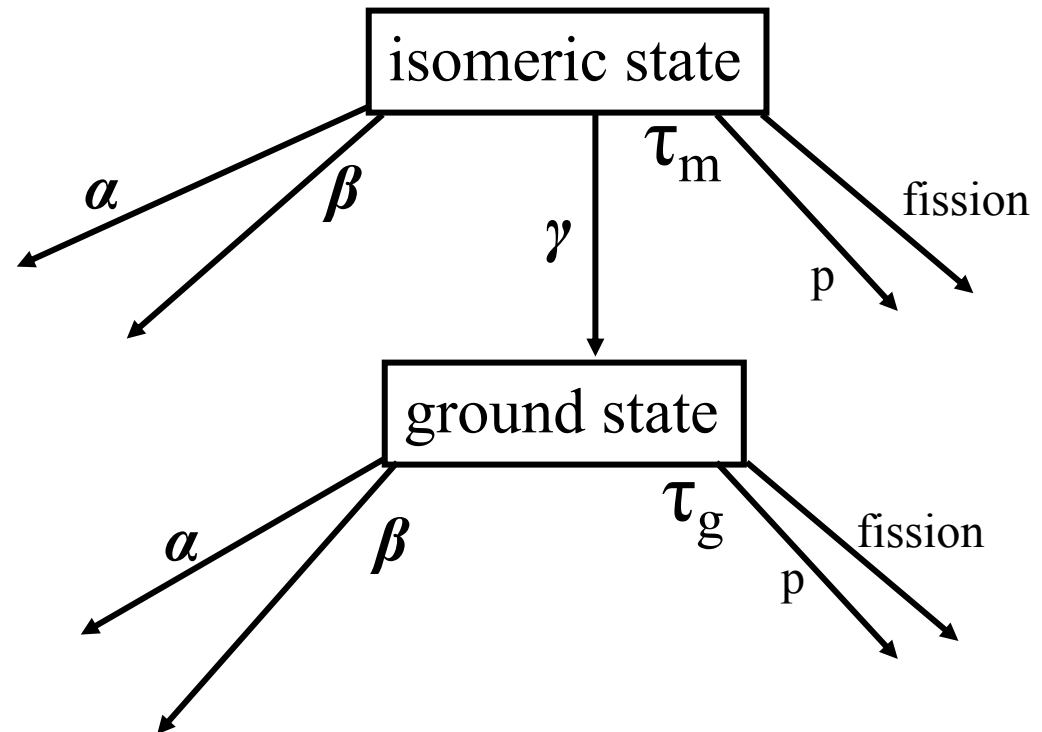
# Nuclear isomers



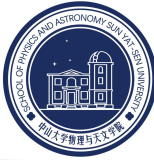
“The existence of isotopic isobars (same-Z, same-A), with clearly distinguishable properties such as different radioactive half-periods, was anticipated in 1917 when Soddy proposed that such nuclei be called isomers if and when found.” **Evans, 1955**

“An excited nuclear state which endures long enough to have a directly measurable lifetime is called an isomeric state.” **Bethe, 1956**

Rule of thumb:  
 $\tau > 1 \text{ ns}$

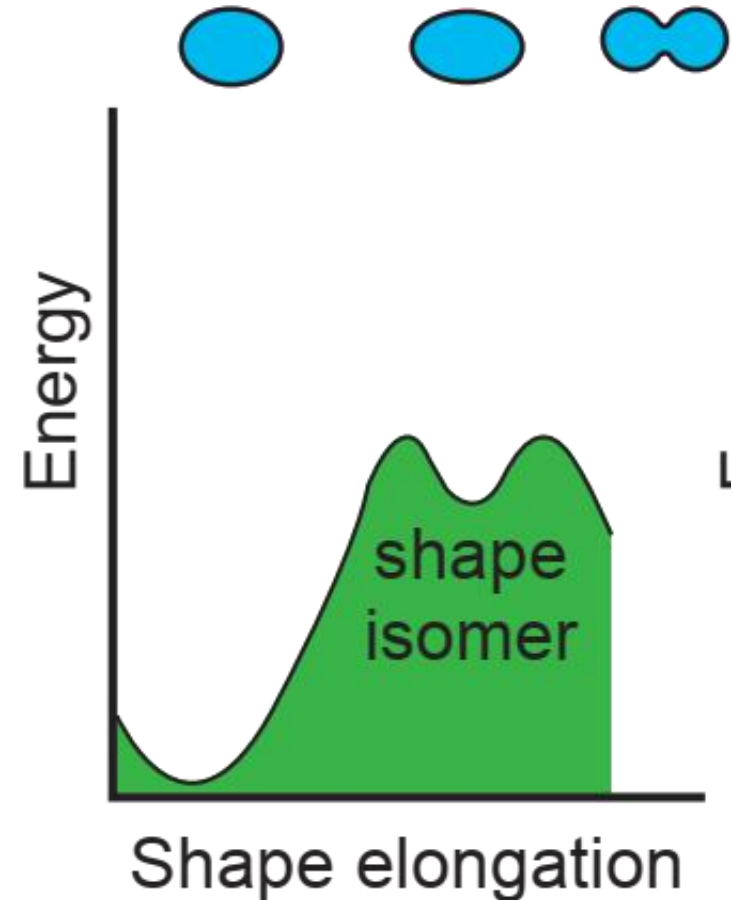


# Types of isomers

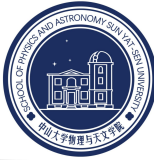


## Shape isomers

- A secondary energy minimum at large elongation.
- Fission isomers in heavy nuclei, such as  $^{242}\text{Am}$ .
- Elongated such that its major-to-minor axis ratio is about 2:1
- Decay back to the ground state competes with fission.



# Types of isomers




## Spin trap

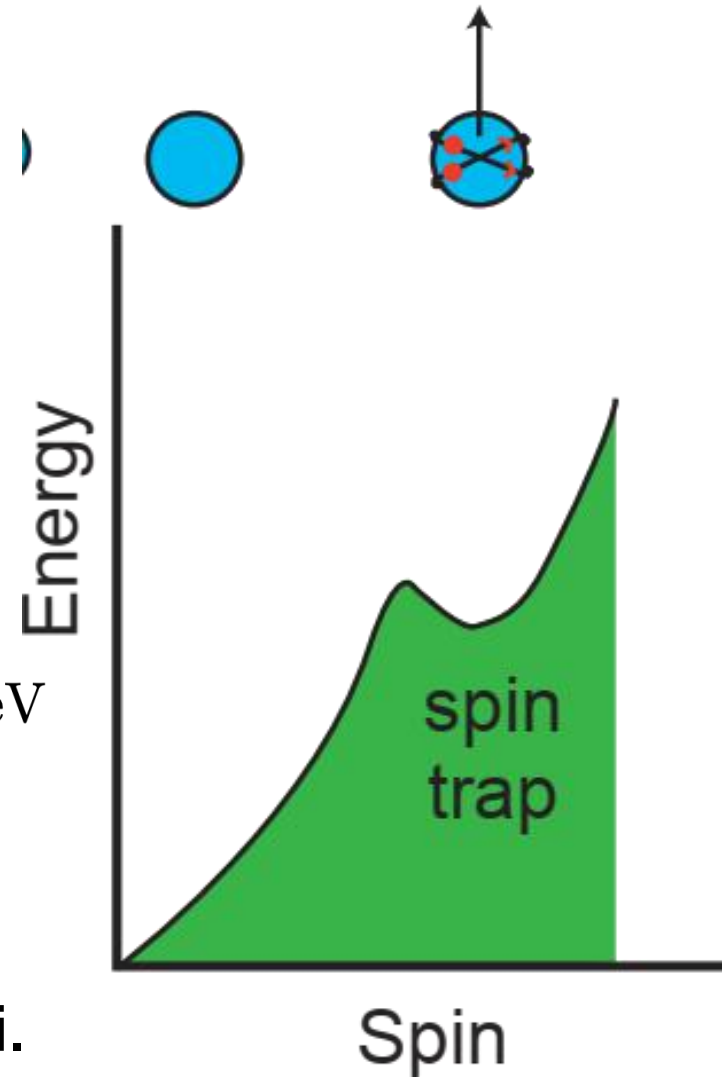
- The decay path to lower energy states requires a large change in nuclear spin, and therefore the emission of radiation with a high multipolarity  $\lambda = \Delta I$

- $^{180m}\text{Ta}$  case:

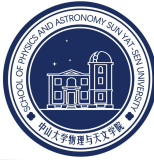
g.s.:  $1^+$ , isomer:  $9^-$ ,  $\lambda = 8$ ,  $E_x = 77.2$  keV

  $T_{1/2}(\text{g.s.}) = 8.154$  h  
 $T_{1/2}(9^-) > 7.1 \times 10^{15}$  a

- In spherical or near-spherical nuclei.



# Types of isomers



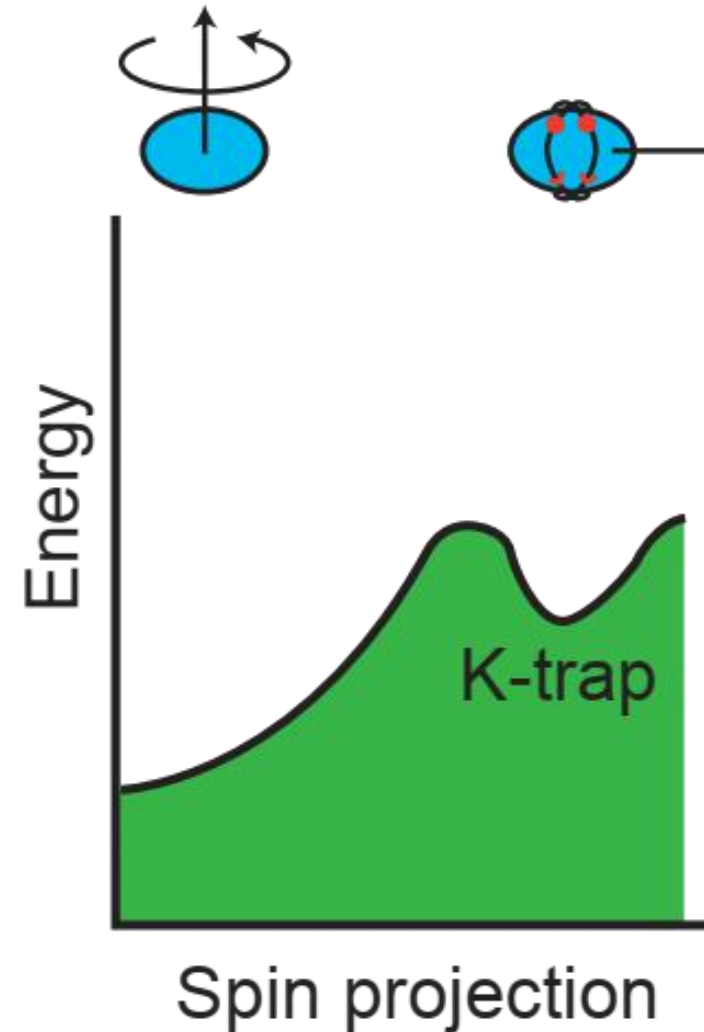
## K traps (K isomers)

- A special spin trap depends on spin's orientation.
- K represents the projection of the total nuclear spin along the symmetric axis of the nucleus.
- $^{180\text{m}}\text{Hf}$  case:  
isomer:  $I = 8, K = 8$ , at 1.1 MeV  
probable decay to  $I = 8, K = 0$  state

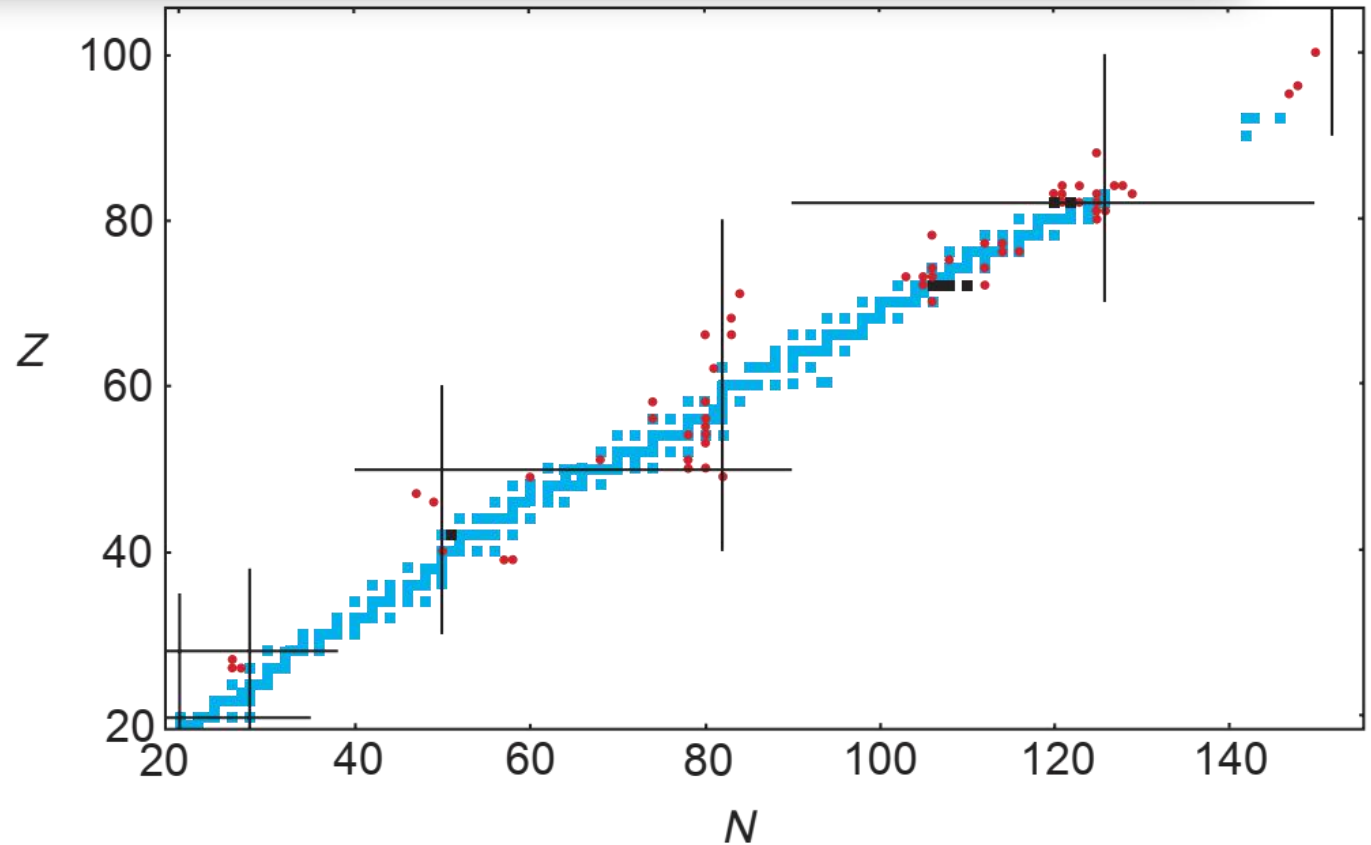
$$\Delta K = 8, \lambda = 0, f = \Delta K - \lambda = 8$$

➡  $T_{1/2}(8^-) = 5.5 \text{ h}$

- In axially deformed (mainly prolate) nuclei.



# Motivation

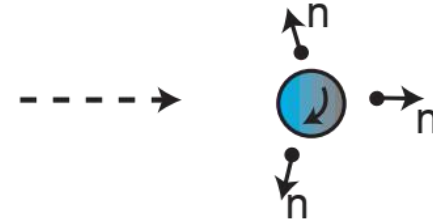
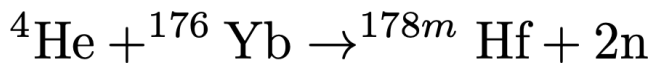


**Figure 2** Chart of atomic nuclei, as a function of neutron number ( $N$ ) and proton number ( $Z$ ). Naturally occurring isotopes are represented by blue squares. The straight lines are drawn for the 'magic numbers' that correspond to closed shells (and spherical shapes). The other symbols indicate isomers with excitation energies greater than 1 MeV and long half-lives: red circles for  $>1$  ms, and black squares for  $>1$  hour. The data are taken from refs. 4, 10, 11.

# Making spin traps



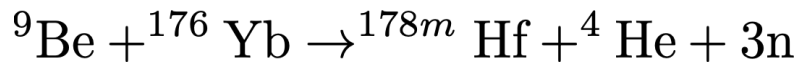
**Fusion:** Colliding  $^4\text{He}$  into  $^{176}\text{Yb}$ , and evaporates two neutrons.



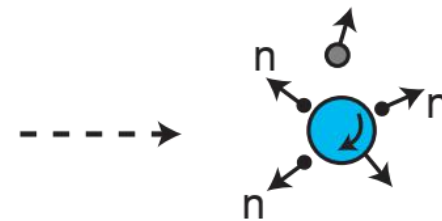
**a**  
Fusion

*bring up to 16 hbar spins and form 31-year  $16^+$  isomer state.*

**Partial fusion:**



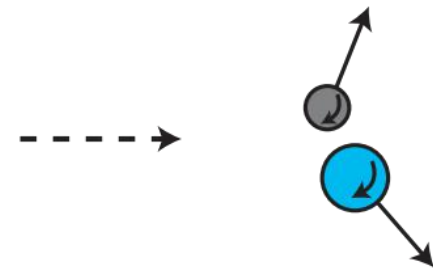
*bring up to 23 hbar spins and form rotational band built on 31-year  $16^+$  isomer state.*



**b**  
Partial fusion

**Scattering reaction:**

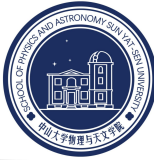
Bombarding heavy, deformed, target nuclei, such as  $^{180}\text{Hf}$  with heavy, deformed, projectile nuclei, such as  $^{238}\text{U}$



**c**  
Scattering

*Nucleons can also be exchanged, so that states in adjacent nuclei may also be reached.*

# Motivation



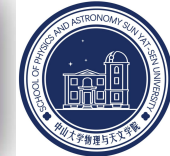
**The purpose of this Paper:** presents new information on the shape evolution of the very neutron-rich  $^{92,94}\text{Se}$  nuclei from an isomer-decay spectroscopy experiment.

## Highlights:

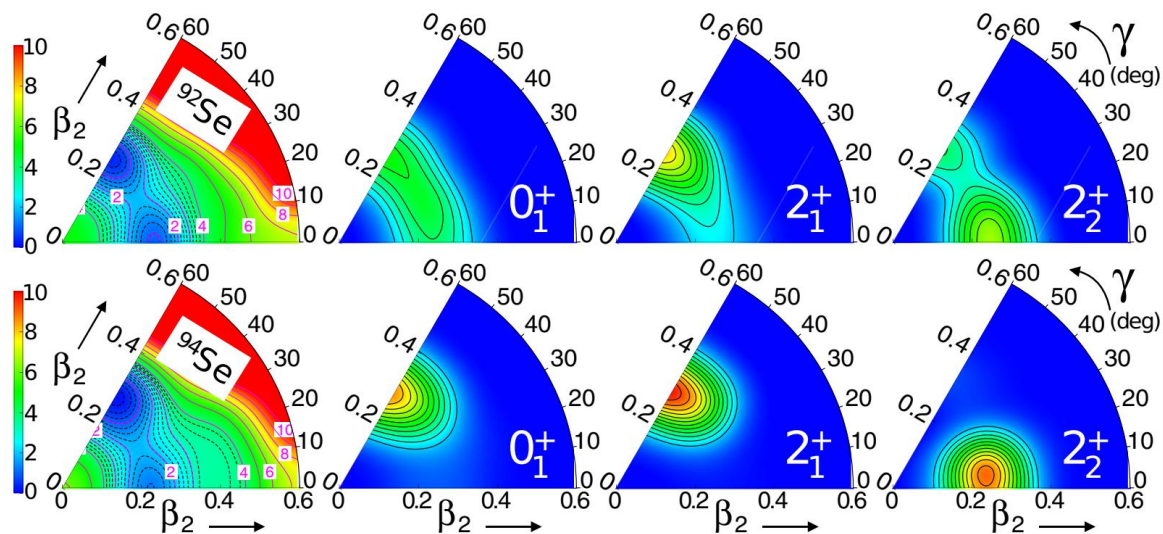
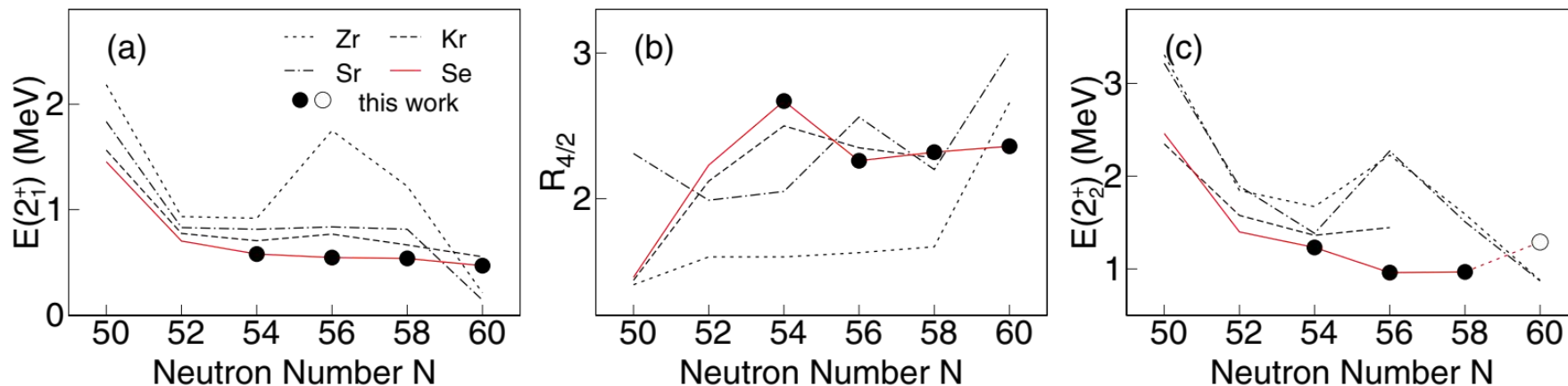
- The isomeric levels are interpreted as originating from high- $K$  quasineutron states with an oblate deformation of  $\beta \sim 0.25$ .
- $^{94}\text{Se}$  is the **lowest-mass** neutron-rich nucleus known to date with such a substantial  $K$  hindrance.
- The first observation of an **oblate  $K$**  isomer in a deformed nucleus.



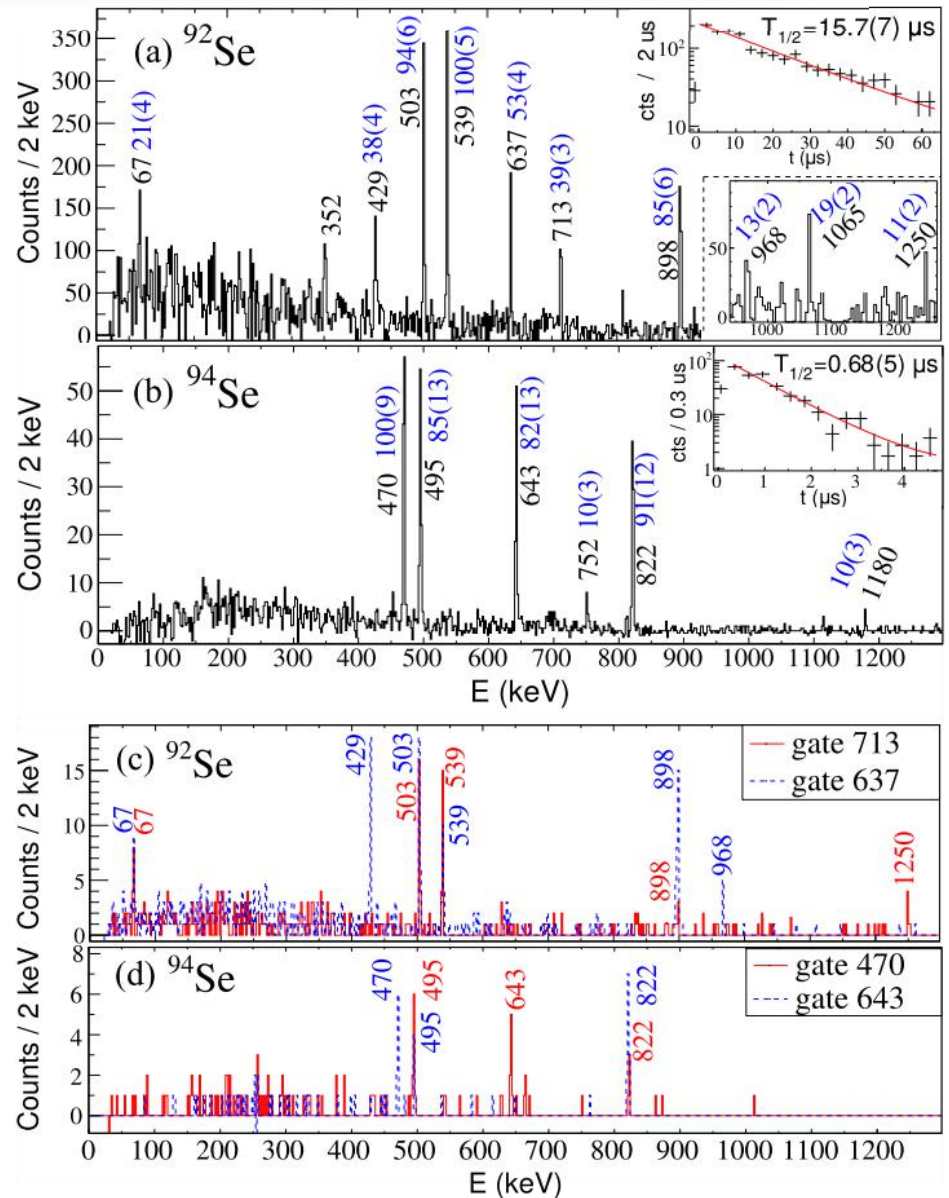
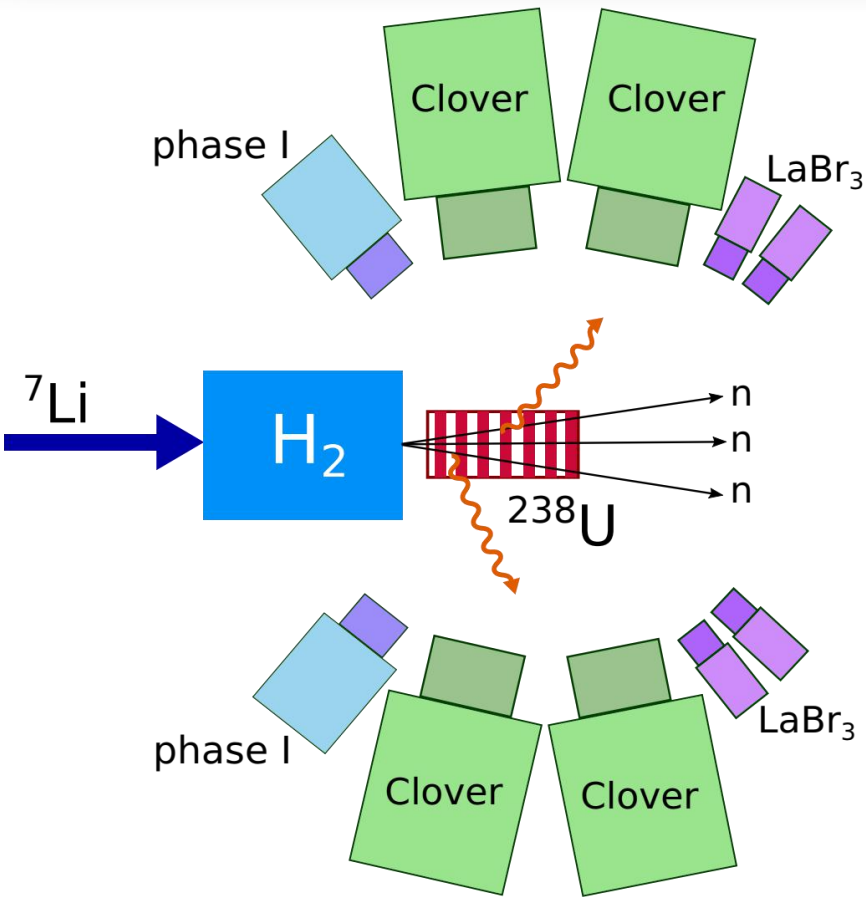
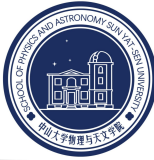
# Motivation



Another thing: both triaxial degree of freedom and shape coexistence playing important roles in the description of intrinsic deformations in neutron-rich Se isotopes.



# Experiment



# Data analysis



For  $^{92}\text{Se}$ ,  $\gamma$  rays detected within  $0.19 \leq t_\gamma \leq 70 \mu\text{s}$  were used.

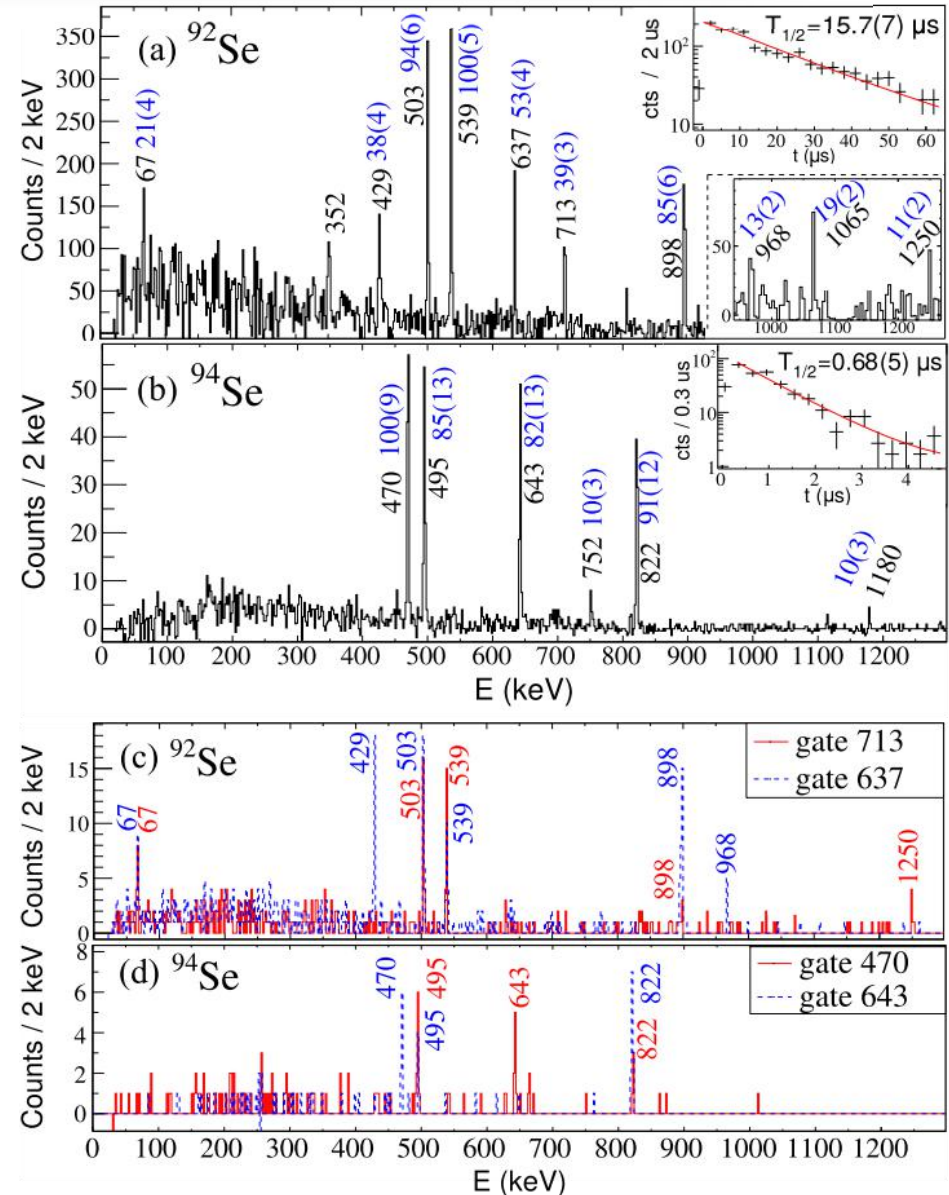
A spectrum of events with  $t_\gamma \geq 70 \mu\text{s}$  was scaled and subtracted to further reduce the environmental background component.

The half-life was determined using an exponential decay curve plus a constant baseline.

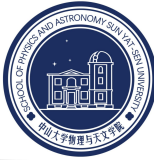
**The half-life of  $K\pi = 7^-$  isomers:**

$$T_{1/2} = 15.7(7) \mu\text{s} \text{ in } ^{92}\text{Se}$$

$$T_{1/2} = 0.68(5) \mu\text{s} \text{ in } ^{94}\text{Se}$$



# Data analysis



## Level schemes constructed for $^{92,94}\text{Se}$

Deduced from energy-sum checks,  $\gamma\gamma$  coincidences, and sums of efficiency-corrected intensities.

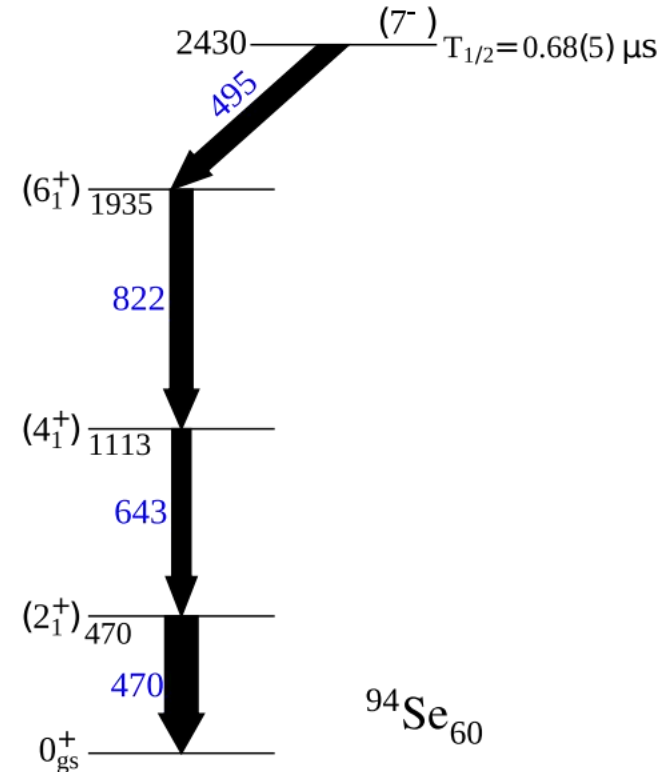
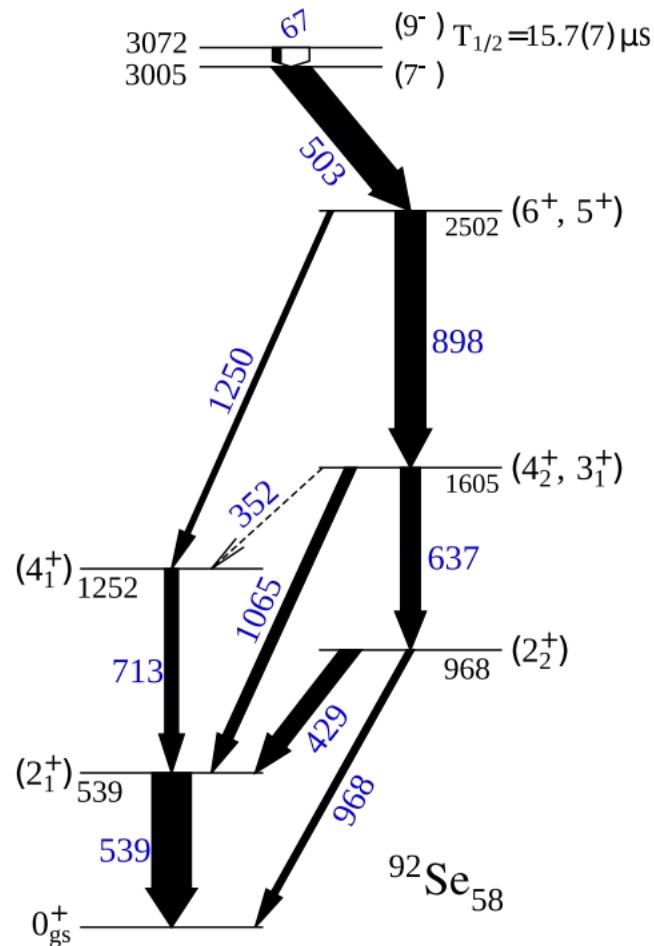
The average internal conversion coefficient of the 67-keV transition is  $\alpha = 3.4(6)$



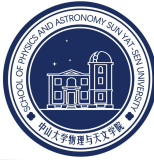
**favors E2 transition**



The 3005- and 3072-keV levels have same parity and  $\Delta J = 2$

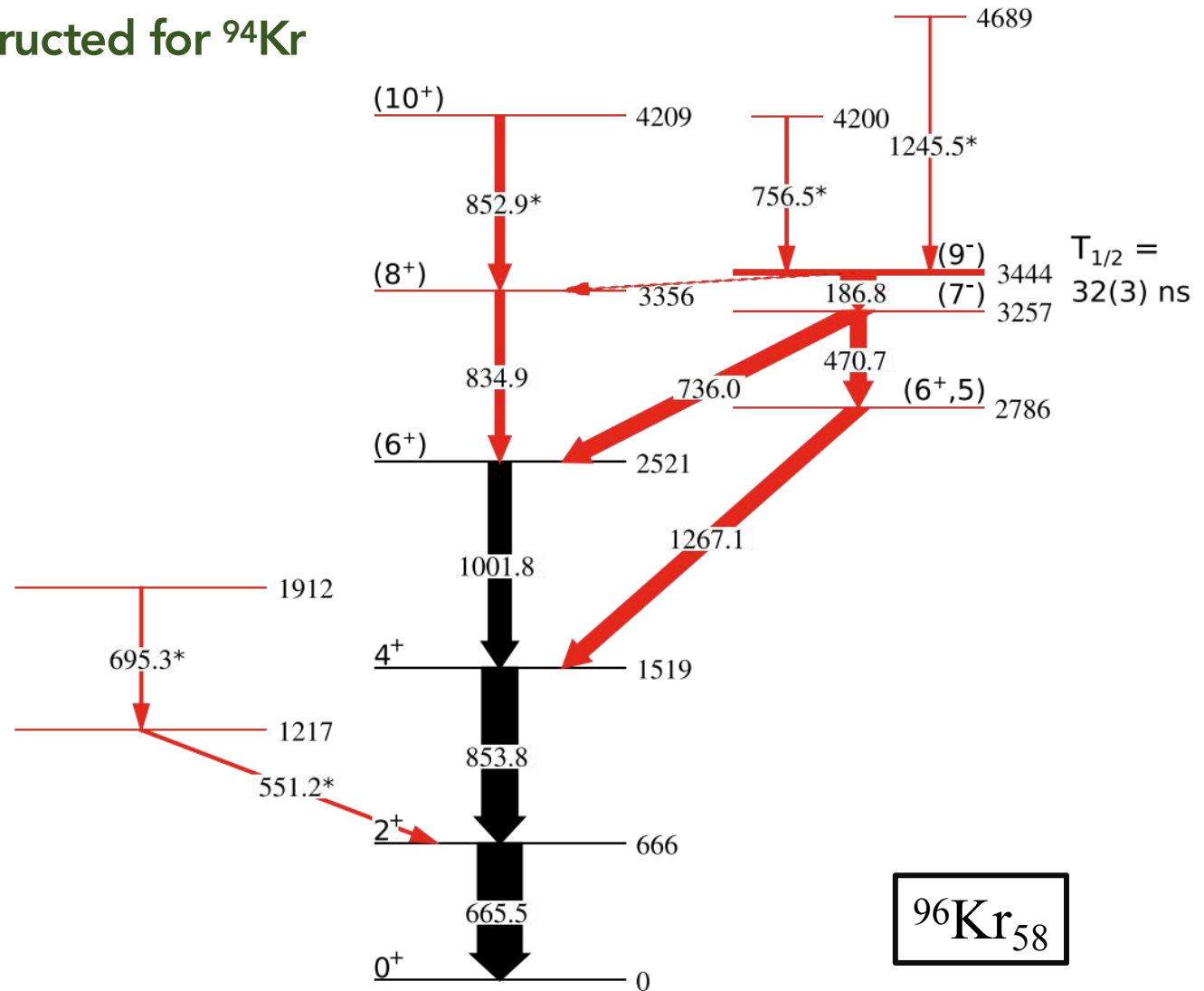


# Data analysis



R. -B. Gerst et. al., PRC 102, 064323 (2020).

## Level schemes constructed for $^{94}\text{Kr}$



# Why $K$ isomers are usually prolate



**Example:** the configuration of high- $K$  isomers in  $^{178}\text{W}$ .

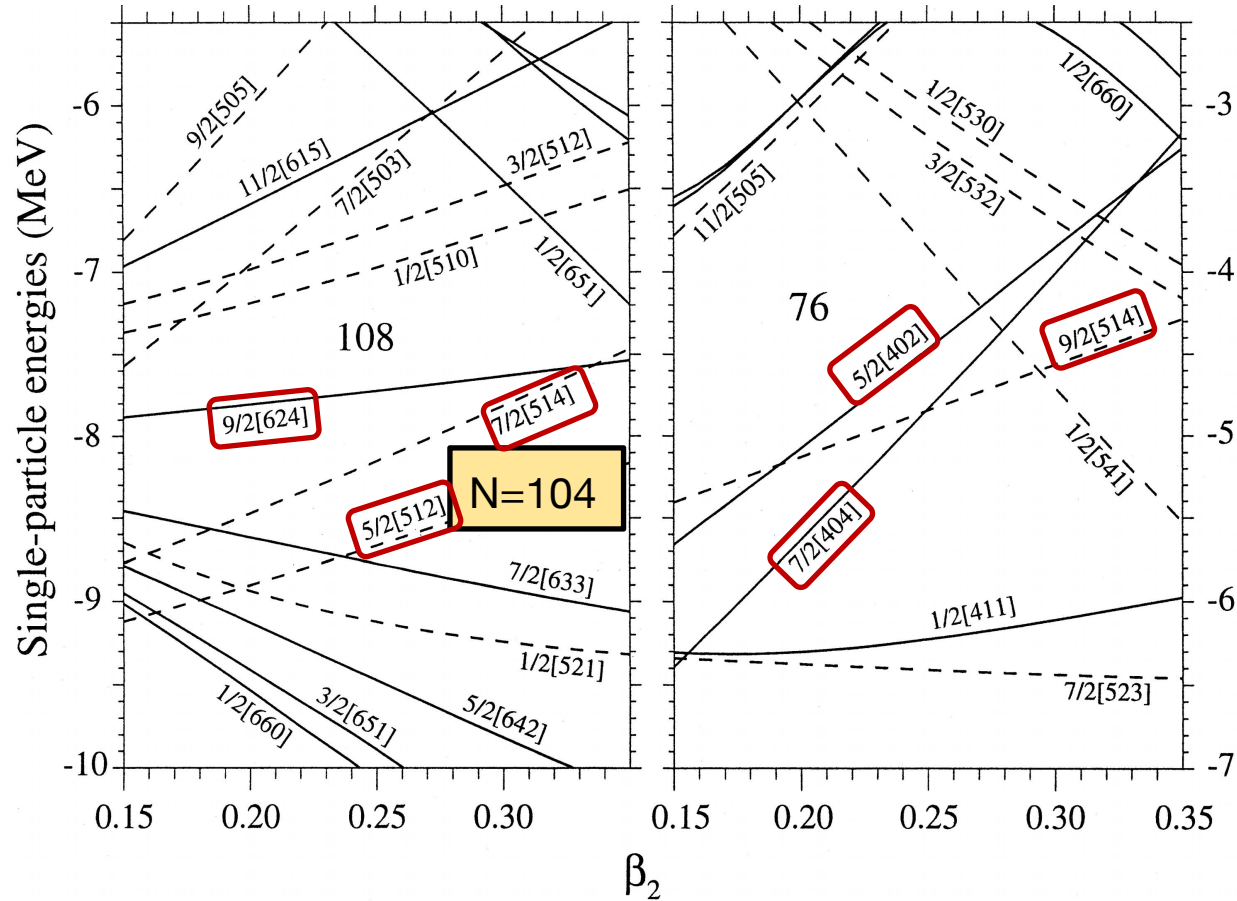
$$K = \Omega_1 + \Omega_2$$

Sum of the projection of angular momentum of the broken-pair nucleons

Large  $K \rightarrow K$  forbidden metastable isomers

$$K\pi = 6^+ : \nu \frac{7}{2} [514] \otimes \nu \frac{5}{2} [512]$$

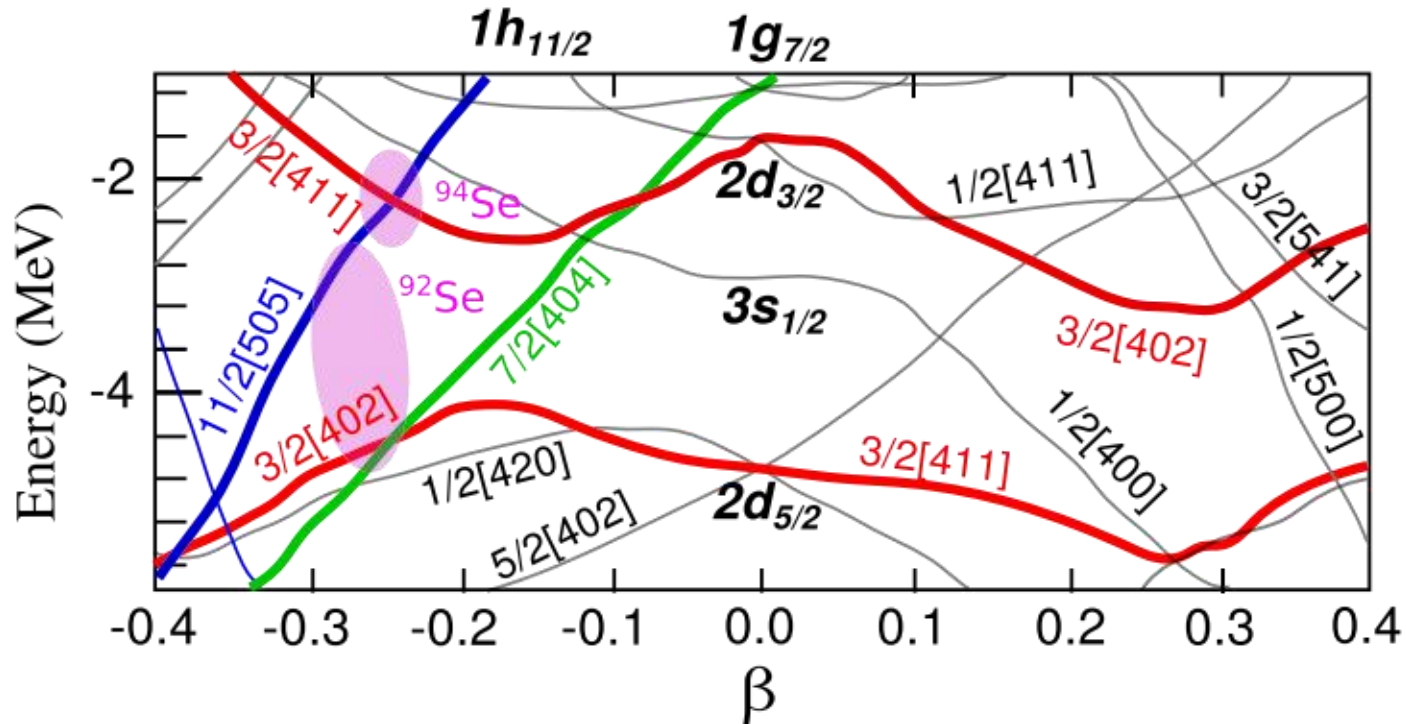
$$K\pi = 8^- : \nu \frac{9}{2} [624] \otimes \nu \frac{7}{2} [514]$$



# Discussions



The Nilsson diagrams obtained by Gogny HFB calculations

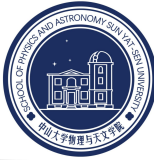


$$K^\pi = 7^-(\nu 11/2^- [505] \otimes \nu 3/2^+ [402])$$

$$K^\pi = 9^-(\nu 11/2^- [505] \otimes \nu 7/2^+ [404])$$

at oblate deformations of  $\beta \sim -0.24$

# Discussions



## Configuration-constrained PES calculations for high- $K$ states

Solve the Schrödinger equation  $H_{\text{WS}} \Psi = E_{\text{WS}} \Psi$ ,

with the Woods-Saxon potential,

$$H_{\text{WS}} = T + V(r, \hat{\beta}) + V_{l \cdot s}(r, \hat{\beta}) + \frac{1}{2} (1 + \tau_3) V_c(r, \hat{\beta}),$$

where

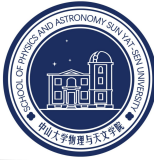
$$V(r, \hat{\beta}) = \frac{-V_{00} \left[ 1 \pm \kappa \frac{N-Z}{N+Z} \right]}{1 + e^{\frac{r-R}{a}}} \quad \text{Central force}$$

$$V_{l \cdot s}(r, \hat{\beta}) = \lambda \left( \frac{\hbar}{2Mc} \right)^2 [\nabla V(r, \hat{\beta}) \times p] \cdot \sigma, \quad \text{Spin-orbit term}$$

$$V_c(r, \hat{\beta}) = \frac{Z-1}{\frac{4}{3}\pi R^3} \int \frac{1}{|r' - r|} d^3 r' \quad \text{Coulomb term}$$



# Discussions



## Configuration-constrained PES calculations for high- $K$ states

Pairing treatment: **Lipkin-Nogami method**

Add a residual pairing interaction to the Hamiltonian

$$H = H_{sp} - G \sum_{k, k' > 0} a_k^\dagger a_{\bar{k}}^\dagger a_{k'} a_{\bar{k}'}$$

And the BCS wave function is

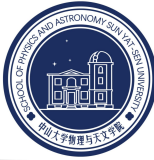
$$|\text{BCS}\rangle = \prod_{k > 0} \left( u_k + v_k a_k^\dagger a_{\bar{k}}^\dagger \right) |0\rangle$$

The  $v_k^2$  and  $u_k^2$  represent the probability that a certain pair state is or is not occupied, which has to be determined in such a way that the corresponding energy has a minimum (variational principle).

The should fulfill  $2 \sum_k v_k^2 = N$

The variational Hamiltonian is  $H' = H - \lambda N$

# Discussions



But BCS method would collapse if the pairing is too weak. Now we add high order particle-number constraint

$$\hat{H}_{LN} = \hat{H} - \lambda \hat{N} - \lambda_2 \hat{N}^2$$

And solve equations iteratively:

$$N = 2 \sum_k v_k^2$$

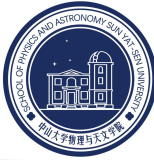
$$\frac{2}{G} = \sum_k \frac{1}{\sqrt{(\epsilon_k - \lambda)^2 + \Delta^2}}$$

$$\epsilon_k = e_k + (4\lambda_2 - G) v_k^2$$

$$v_k^2 = \frac{1}{2} \left[ 1 - \frac{\epsilon_k - \lambda}{\sqrt{(\epsilon_k - \lambda)^2 + \Delta^2}} \right]$$

$$\lambda_2 = \frac{G}{4} \left[ \frac{(\sum_k u_k^3 v_k) (\sum_k u_k v_k^3) - \sum_k u_k^4 v_k^4}{(\sum_k u_k^2 v_k^2)^2 - \sum_k u_k^4 v_k^4} \right]$$

# Discussions



The expectation value of the total Hamiltonian becomes

$$E_{LN} = \sum_k 2v_k^2 e_k - \frac{\Delta^2}{G} - G \sum_k V_k^4 - 4\lambda_2 \sum_k (u_k v_k)^2$$

If we want to create a broken pair, block the  $i$ - and  $j$ -th single-particle orbitals, make sure that in every iteration the  $v_i = 1$ ,  $v_j = 1$ . And the  $LN$  energy is

$$E_{LN} = \sum_{j=1}^S e_{k_j} + \sum_{k \neq k_j} 2v_k^2 e_k - \frac{\Delta^2}{G} - G \sum_{k \neq k_j} v_k^4 + G \frac{N-S}{2} - 4\lambda_2 \sum_{k \neq k_j} (u_k v_k)^2.$$

Using macroscopic-microscopic method, i.e.,

$$E_{\text{tot}} = E_{\text{LDM}} + E_{\text{LN}} - \tilde{E}_{\text{strut}}.$$

Calculate the energies with different  $\beta_2$ ,  $\beta_4$ , and  $\gamma$  to get the configuration-constrained PES.

# Discussions



## Results given by Configuration-constrained PES

For  $^{92}\text{Se}$ : g.s. deformation is **prolate**

$$K^\pi = 7^-(\nu 11/2^- [505] \otimes \nu 3/2^+ [402])$$

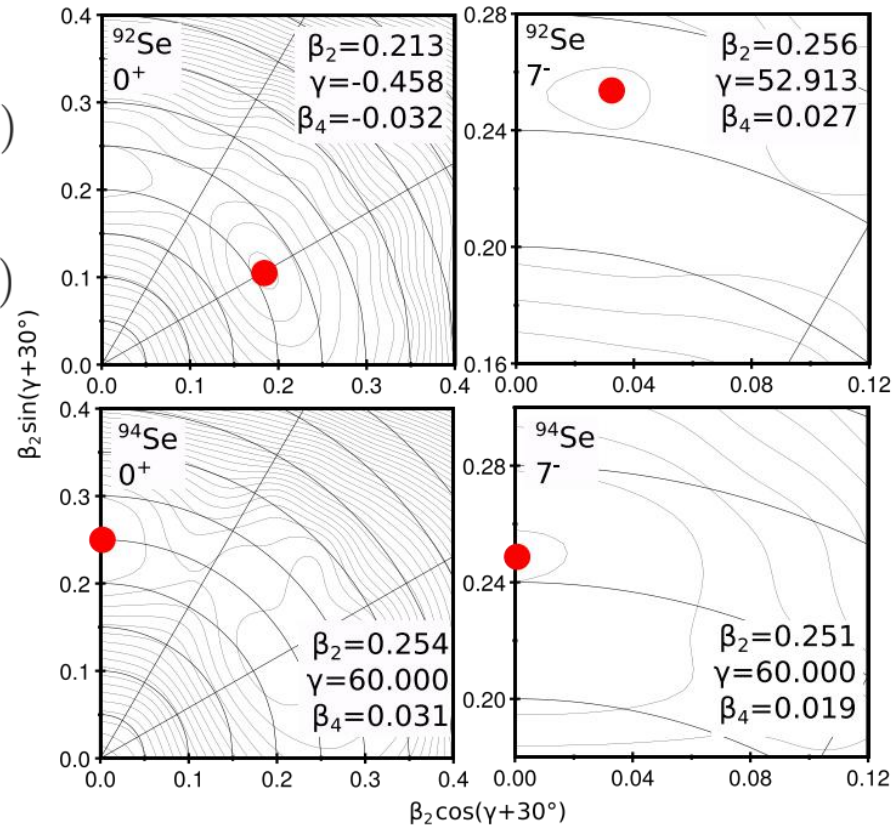
deformation:  $(\beta_2, \gamma, \beta_4) = (0.256, 52.9, 0.027)$

$$K^\pi = 9^-(\nu 11/2^- [505] \otimes \nu 7/2^+ [404])$$

deformation:  $(\beta_2, \gamma, \beta_4) = (0.254, 54.1, 0.021)$

表 1  $N=58$  同中子链中  $K^\pi = 9^-$  态激发能与形变参数的计算结果\*

核素	$\beta_2$ (基态)	$\beta_2$ (激发态)	$\gamma$ (基态) / $^\circ$	$\gamma$ (激发态) / $^\circ$	$E_{\text{cal}}$ /keV	$E_{\text{exp}}$ /keV
$^{88}\text{Zn}$	0.182	0.183	0	45	6 026	
$^{90}\text{Ge}$	0.198	0.196	1	58	3 286	
$^{92}\text{Se}$	0.208	0.209	0	60	3 255	3 072
$^{94}\text{Kr}$	0.253	0.238	0	58	3 322	3 444
$^{96}\text{Sr}$	0.315	0.206	0	60	3 853	3 524
$^{98}\text{Zr}$	0.309	0.199	0	57	3 522	
$^{100}\text{Mo}$	0.202	0.195	35	54	3 362	
$^{102}\text{Ru}$	0.195	0.180	24	0	3 353	
$^{104}\text{Pd}$	0.140	0.164	0	20	4 987	
$^{106}\text{Cd}$	0.127	0.127	0	15	5 809	



# Discussions



## Results given by Configuration-constrained PES

For  $^{94}\text{Se}$ : g.s. deformation is **oblate**

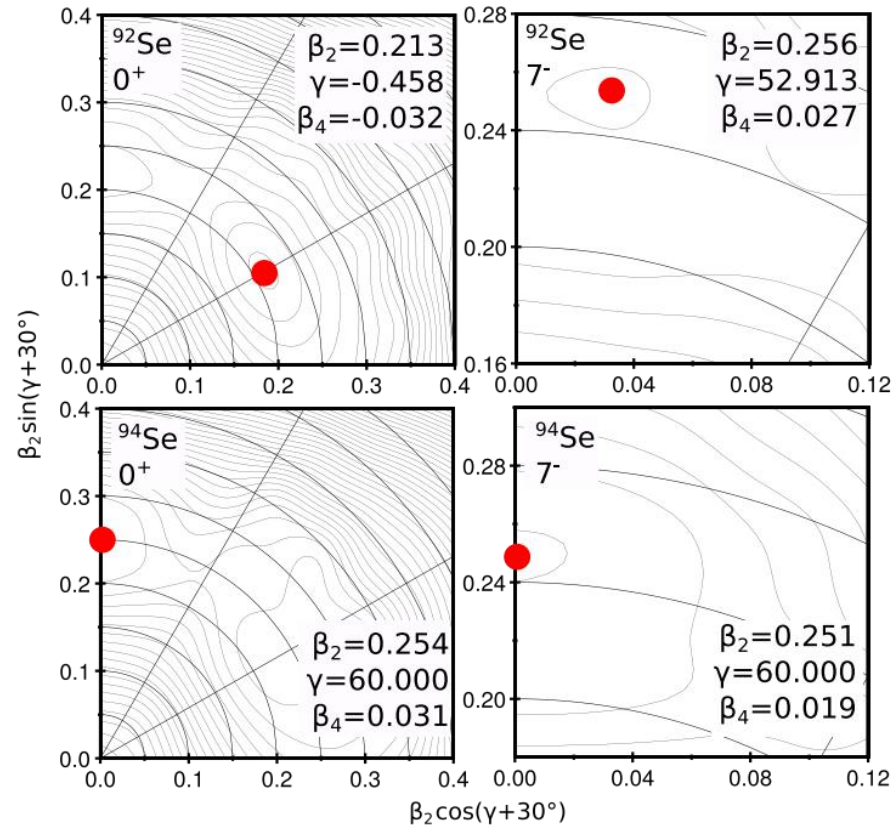
$$K^\pi = 7^-(\nu 11/2^- [505] \otimes \nu 3/2^+ [402])$$

deformation:

$$(\beta_2, \gamma, \beta_4) = (0.251, 60.0, 0.019)$$

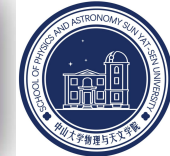
表2  $N=60$  同中子链中  $K^\pi = 7^-$  态激发能与形变参数的计算结果\*

核素	$\beta_2$ (基态)	$\beta_2$ (激发态)	$\gamma$ (基态) /( $^\circ$ )	$\gamma$ (激发态) /( $^\circ$ )	$E_{\text{cal}}$ /keV	$E_{\text{exp}}$ /keV
$^{90}\text{Zn}$	0.195	0.207	56	47	2 259	
$^{92}\text{Ge}$	0.210	0.219	1	41	2 353	
$^{94}\text{Se}$	0.239	0.236	58	58	2 431	2 400
$^{96}\text{Kr}$	0.309	0.254	0	59	2 837	
$^{98}\text{Sr}$	0.334	0.221	0	57	3 791	
$^{100}\text{Zr}$	0.340	0.213	1	50	3 511	
$^{102}\text{Mo}$	0.262	0.262	21	15	2 578	
$^{104}\text{Ru}$	0.226	0.233	31	20	2 621	
$^{106}\text{Pd}$	0.174	0.188	0	2	1 786	
$^{108}\text{Cd}$	0.134	0.159	0	1	2 043	

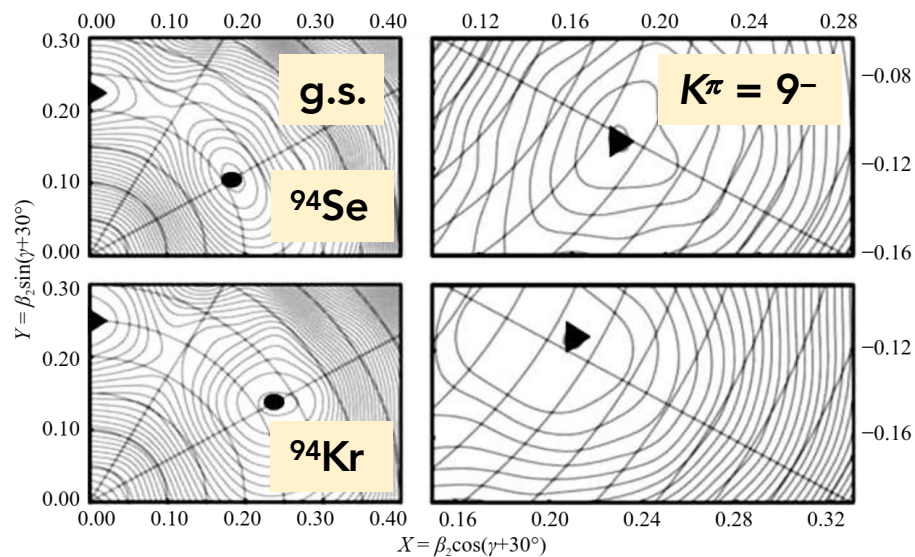


The rather soft potential energy surface implying that  $K$  is a less well-defined quantum number.

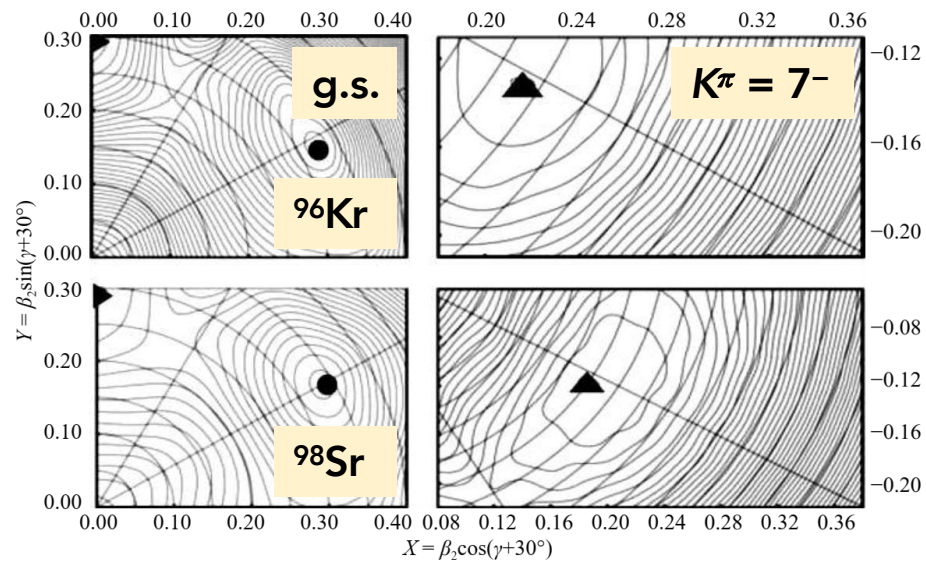
# Discussions



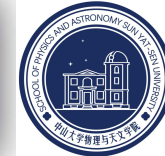
## More results given by Configuration-constrained PES



## Shape + K isomer?

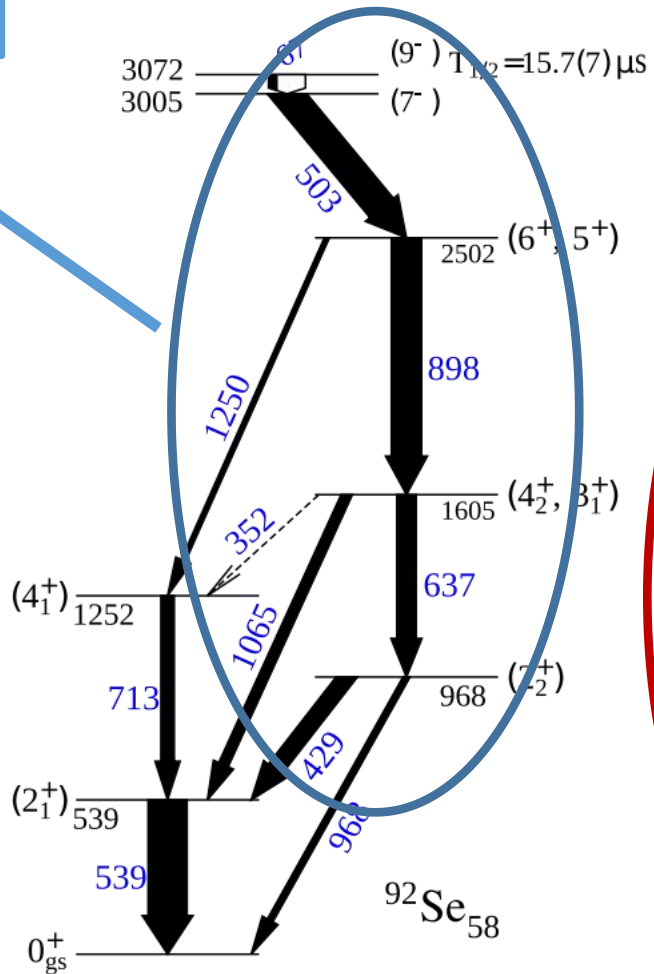
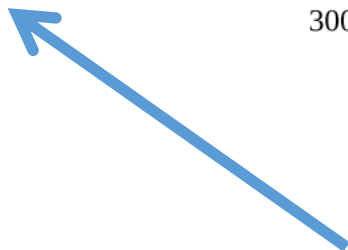


# Discussions



A question that hasn't been answered

Decay to yrare bands



Shape +  $K$  isomer?

Decay to yrast bands

