



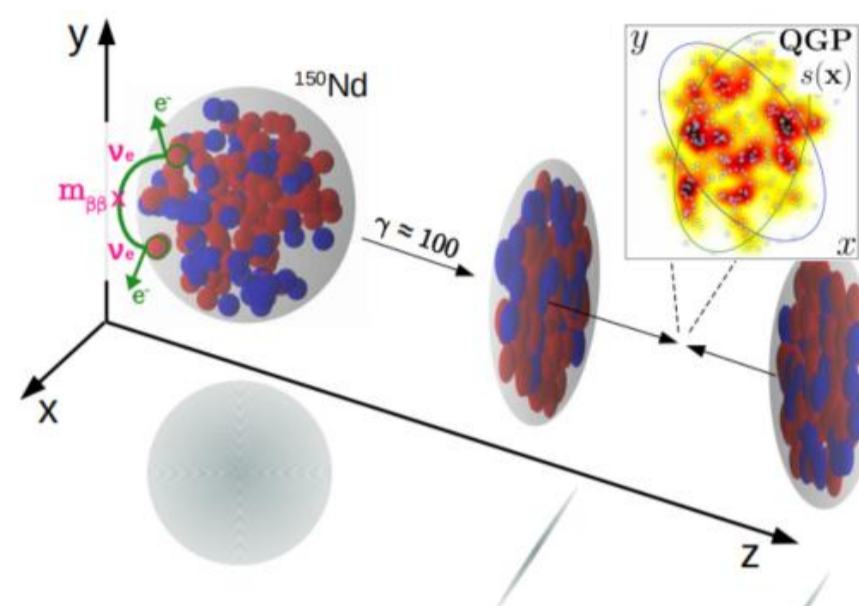
Benchmarking nuclear matrix elements of $0\nu\beta\beta$ decay with high-energy nuclear collisions



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Introduction

- Neutrinoless double-beta decay ($0\nu\beta\beta$) plays a crucial role in determining whether neutrinos are Dirac or Majorana particles and in shedding light on the origin of matter-antimatter asymmetry in the universe.
- The nuclear matrix elements (NMEs) connect the observed half-lives of $0\nu\beta\beta$ decay to the underlying neutrino masses; however, predictions from different nuclear models can vary by factors of three or more, leading to significant theoretical uncertainties. Reducing these uncertainties remains a central challenge in both the design and interpretation of experiments aimed at discovering $0\nu\beta\beta$ decay.
- Exploring correlations between NMEs and experimentally accessible observables offers a promising approach to constraining and reducing these uncertainties.
- This work: we study the correlation between the NME of $0\nu\beta\beta$ decay $^{150}\text{Nd} \rightarrow ^{150}\text{Sm}$ and the geometric properties of the quark-gluon plasma (QGP) that are produced in the ultra-central relativistic heavy-ion collision (HIC) of $^{150}\text{Nd} + ^{150}\text{Nd}$.



Methods

Relativistic energy density functional (EDF)

- The EDF is composed of kinetic energy, electromagnetic energy and interaction energy

$$E[\tau, \rho, \nabla\rho; \mathbf{C}] = \int d^3r [\tau(\mathbf{r}) + \mathcal{E}^{\text{em}}(\mathbf{r}) + \mathcal{E}^{\text{int}}(\mathbf{r})].$$

- The NN interaction EDF

$$\begin{aligned} \mathcal{E}^{\text{int}}(\mathbf{r}) = & \frac{\alpha_S}{2} \rho_S^2 + \frac{\beta_S}{3} \rho_S^3 + \frac{\gamma_S}{4} \rho_S^4 + \frac{\delta_S}{2} \rho_S \Delta \rho_S \\ & + \frac{\alpha_V}{2} j_\mu j^\mu + \frac{\gamma_V}{4} (j_\mu j^\mu)^2 + \frac{\delta_V}{2} j_\mu \Delta j^\mu \\ & + \frac{\alpha_{TV}}{2} \mathbf{j}_{TV}^\mu \cdot (\mathbf{j}_{TV})_\mu + \frac{\delta_{TV}}{2} \mathbf{j}_{TV}^\mu \cdot \Delta (\mathbf{j}_{TV})_\mu \end{aligned}$$

- The free parameters

$$\mathbf{C} = \{\alpha_S, \beta_S, \gamma_S, \delta_S, \alpha_V, \gamma_V, \delta_V, \alpha_{TV}, \delta_{TV}\}$$

MR-CDFT for nuclear low-lying states

- Nuclear wave functions are linear combinations of particle-number (N, Z) and angular-momentum(J) projected wave functions,

$$|\Psi_{I/F}(J_V^+, \mathbf{C})\rangle = \sum_q f_v^{JNZ}(\mathbf{q}, \mathbf{C}) \hat{P}^N \hat{P}^Z \hat{P}^J |\Phi(\mathbf{q}, \mathbf{C})\rangle,$$

where weight function f_v^{JNZ} is determined by the variational principles.

- The mean-field wave functions $|\Phi(\mathbf{q}, \mathbf{C})\rangle$ are determined by deformation constrained relativistic mean-field (RMF) + BCS theory.
- The NME of ground state to the ground state $0\nu\beta\beta$ decay is given by

$$M^{0\nu}(\mathbf{C}) = \langle \Psi_F(0_1^+, \mathbf{C}) | \hat{\mathcal{O}}^{0\nu} | \Psi_I(0_1^+, \mathbf{C}) \rangle.$$

Development of SP-CDFT

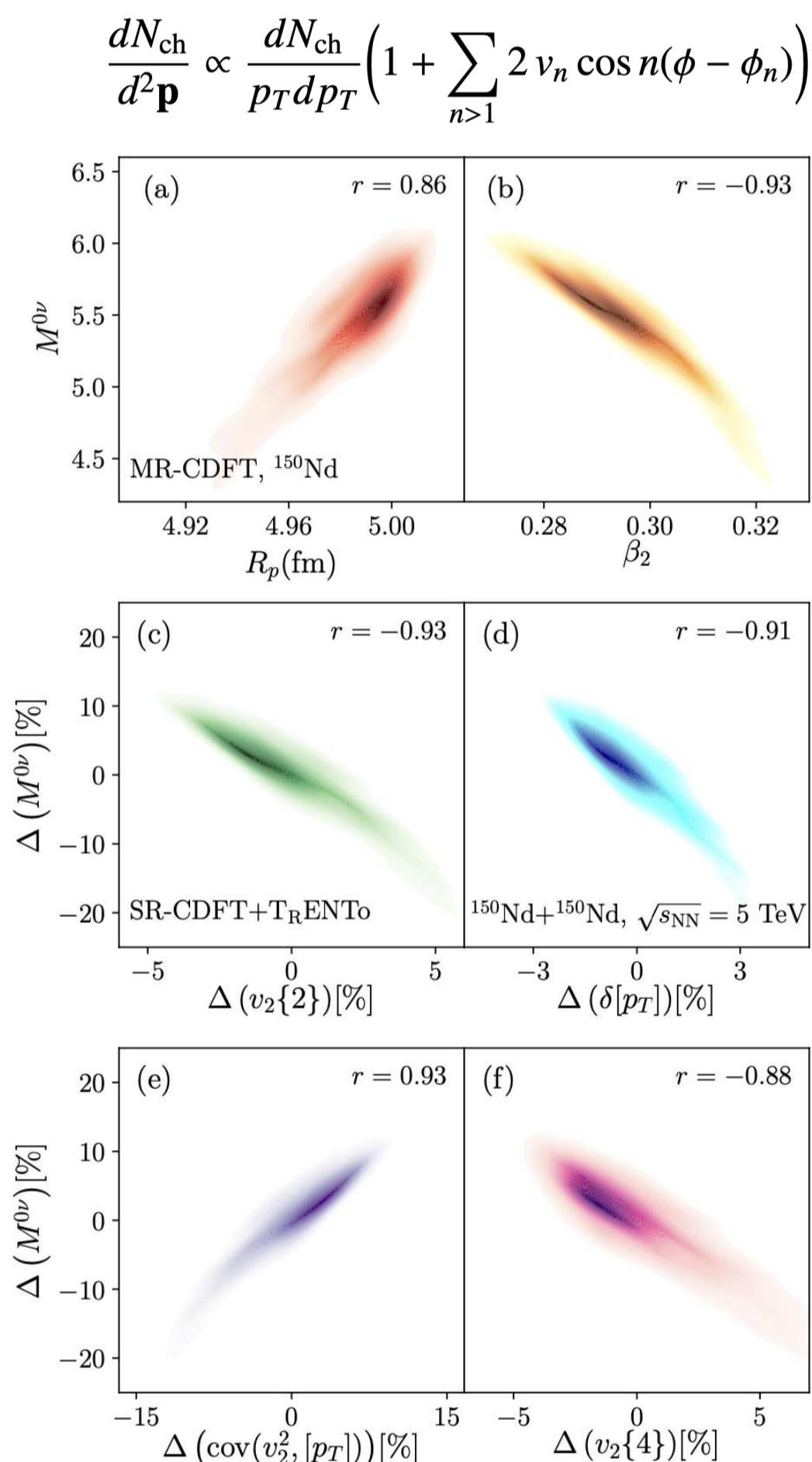
- An emulator based on the eigenvector continuation method, capable of accurately reproducing MR-CDFT results for nuclear low-lying states with significantly reduced computational cost.
- The wave function $|\Psi^{JNZ}(J_k^+, \mathbf{C}_\odot)\rangle$ for the given parameter set \mathbf{C}_\odot is expanded over the MR-CDFT wave functions $|\Psi^{JNZ}(J_v^+, \mathbf{C}_t)\rangle$ of the training set $\{\mathbf{C}_t\}$

$$|\Psi^{JNZ}(J_k^+, \mathbf{C}_\odot)\rangle = \sum_{v=1}^{k_{\max}} \sum_{t=1}^{N_t} f_{k, \mathbf{C}_\odot}^{JNZ}(v, \mathbf{C}_t) |\Psi^{JNZ}(J_v^+, \mathbf{C}_t)\rangle.$$

$T_{\text{R}}\text{ENTo}$ for heavy-ion collisions

- The nuclear density as input for the $T_{\text{R}}\text{ENTo}$ model for the high-energy $^{150}\text{Nd} + ^{150}\text{Nd}$ collisions.
- Computing the spatial ellipticity ε_2 and the ratio of total energy to total entropy E/S of the QGP, ultimately yielding the final-state elliptic flow v_2 and average transverse momentum $[p_T]$ of the collective flow.

Results



- Spreads of the NMEs and the observables in $^{150}\text{Nd} + ^{150}\text{Nd}$ collision by 1k parameter sets, which predict different B(E2) values for ^{150}Nd .
- The existence of strong correlations between the value of the NME of $0\nu\beta\beta$ decay and quantities accessible at colliders.

Summary, Conclusion & Outlook

Summary & Conclusion

- We have conducted the first statistical analysis of the correlation between NMEs of $0\nu\beta\beta$ decay and observables from HIC using covariant density functional theory (CDFT).
- Our study reveals strong correlations between the NMEs and specific HIC observables, whose precise experimental measurement could serve as a valuable benchmark for the prediction of the NMEs in the future.

Outlook

- Extension to other candidate nuclei of $0\nu\beta\beta$ decay.
- Exploration of the correlation between the NMEs with other observables, such as the vector meson production cross section in the HIC.

References:

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- [3] J.M. Yao et al., Prog. Part. Nucl. Phys. 126, 103965 (2022)